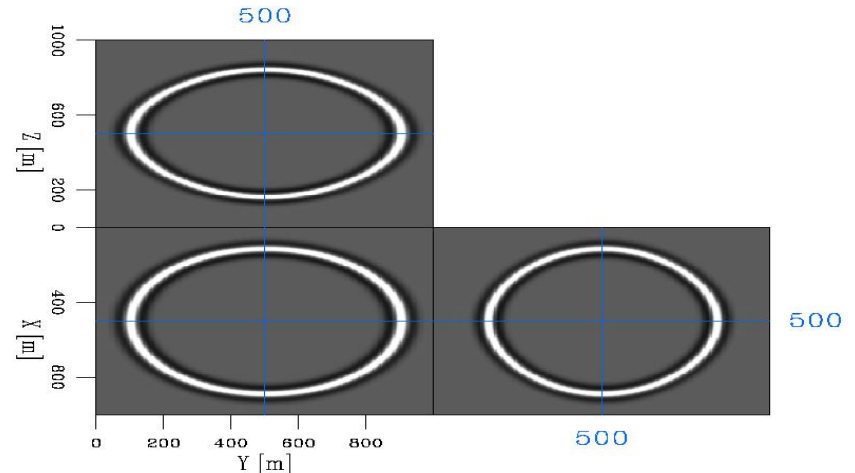
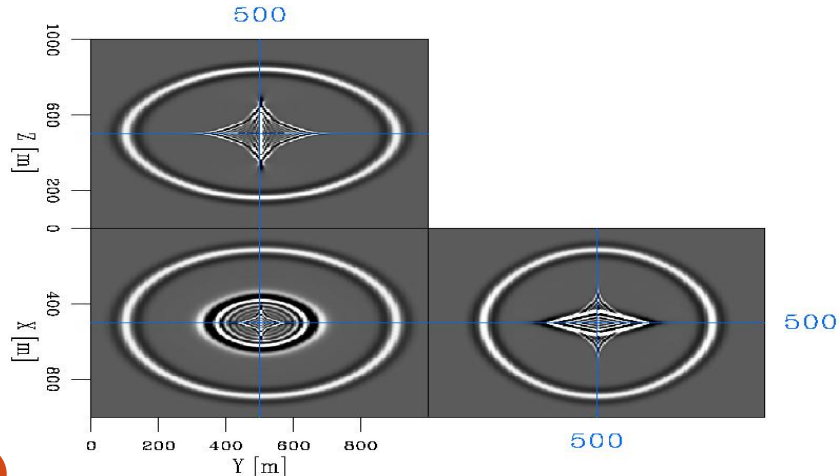


Removing Shear Artifacts in Acoustic Anisotropic Modeling

Huy Le, Stewart A. Levin, Robert G. Clapp, and Biondo Biondi*
SEP 152, p. 129-140



Outline

1. Acoustic approximation and shear artifacts
2. Existing methods
3. Proposed method
4. Accuracy and cost estimates
5. Application to inhomogeneous media
6. Conclusions

Shear artifacts

- Exact wave equations in anisotropic media couple P and S.
- Acoustic approximation: setting S-velocity **along symmetry axes** to 0 (Alkhalifah, 1998)
- S-velocity is **not 0 everywhere** (Grechka et al., 2004).

Existing methods

1. Put the source in isotropic region (Alkhalifah, 1998)
 - Limited application
 - Converted shear artifacts
2. In VTI, use finite shear velocity (Fletcher et al., 2004)
 - What shear velocity to choose?
 - Real shear wave
3. Factorize the dispersion relation (majority)
 - Applicable with source in any media
 - No shear velocity choice

Proposed method

- Factorize dispersion relation to decouple P and S
 - Dispersion relation in VTI:
cubic = quadratic of P and SV + linear of SH
- ⇒ factorization can be done analytically.
- Similar to elastic wave mode separation (Dellinger and Etgen, 1990; Yan and Sava, 2009)
 - In orthorhombic: **eigenvalue decomposition**

Acoustic wave equations in orthorhombic media

$$\frac{\partial^2}{\partial t^2} \boldsymbol{\sigma} = \mathbf{M} \mathbf{D} \boldsymbol{\sigma}$$

$$\boldsymbol{\sigma} = [\sigma_{11} \quad \sigma_{22} \quad \sigma_{33}]^T$$

$$\mathbf{M} = \frac{1}{\rho} \begin{bmatrix} c_{11} & c_{12} & c_{13} \\ c_{12} & c_{22} & c_{23} \\ c_{13} & c_{23} & c_{33} \end{bmatrix}, \mathbf{D} = \begin{bmatrix} \frac{\partial^2}{\partial x^2} & & \\ & \frac{\partial^2}{\partial y^2} & \\ & & \frac{\partial^2}{\partial z^2} \end{bmatrix}$$

Decoupling orthorhombic wave equations

$$\frac{\partial^2}{\partial t^2} \boldsymbol{\sigma} = \mathbf{MD}\boldsymbol{\sigma}$$

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$$\boldsymbol{\sigma}^q = \mathbf{Q}^{-1}\boldsymbol{\sigma} : \text{projection on row vectors}$$

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$$\boldsymbol{\sigma}^q = \mathbf{Q}^{-1}\boldsymbol{\sigma} : \text{projection on row vectors}$$

$$\frac{\partial^2}{\partial t^2} \sigma_i^q = \boldsymbol{\lambda}_i \sigma_i^q$$

$$i = 1, 2, 3$$

Decoupling orthorhombic wave equations

$$\frac{\partial^2}{\partial t^2} \boldsymbol{\sigma} = \mathbf{M}\mathbf{D}\boldsymbol{\sigma}$$

system of 3
coupled equations

$$\mathbf{M}\mathbf{D} = \mathbf{Q}\boldsymbol{\Lambda}\mathbf{Q}^{-1}, \boldsymbol{\Lambda} = \text{diag}[\boldsymbol{\lambda}_i]$$

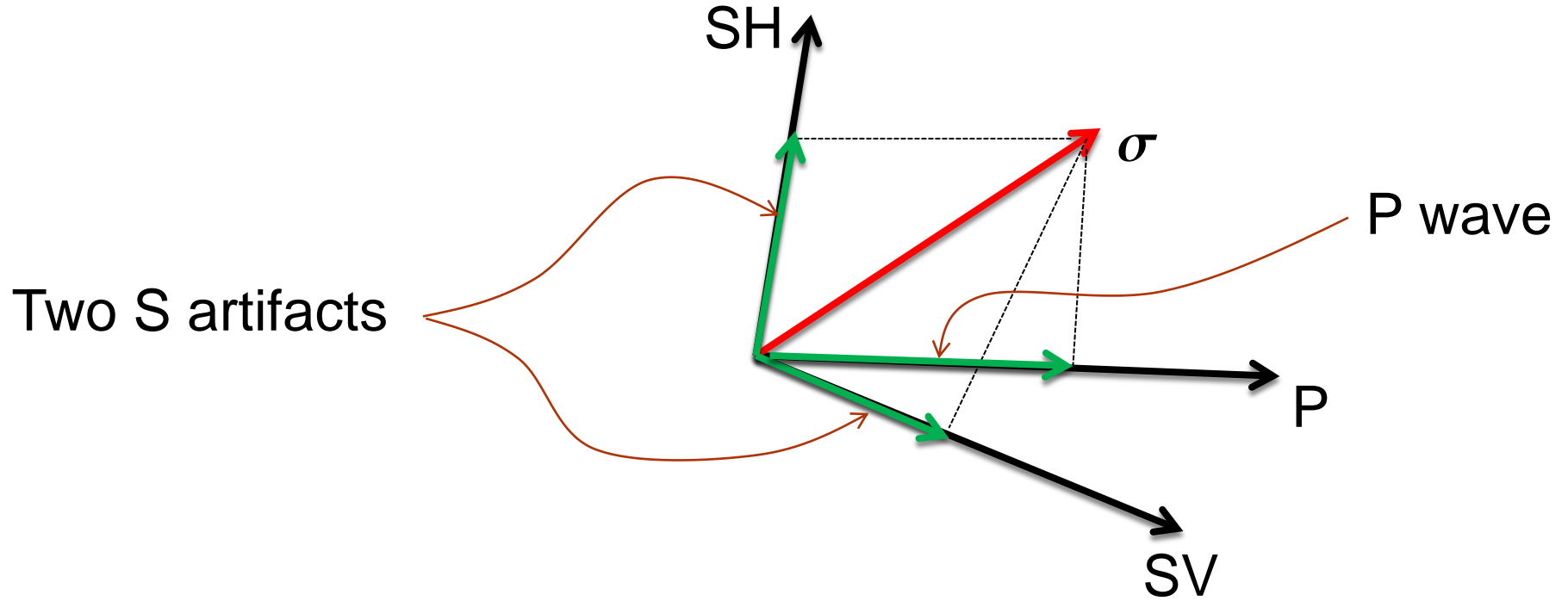
$$\boldsymbol{\sigma}^q = \mathbf{Q}^{-1}\boldsymbol{\sigma} : \text{projection on row vectors}$$

$$\frac{\partial^2}{\partial t^2} \sigma_i^q = \boldsymbol{\lambda}_i \sigma_i^q$$

3 decoupled
equations

$$i = 1, 2, 3$$

Projections on row vectors



Degenerate case: Isotropic

$$\lambda_1 = -k_x^2 - k_y^2 - k_z^2 \quad : \text{Laplacian}$$

$$\lambda_2 = \lambda_3 = 0$$

$$\boldsymbol{\sigma}^q = \begin{bmatrix} p \\ 0 \\ 0 \end{bmatrix} \quad : \text{pressure}$$

Proposed method: use **eigenvalues** as propagators
(elastic wave mode separation: projection on **eigenvectors**)

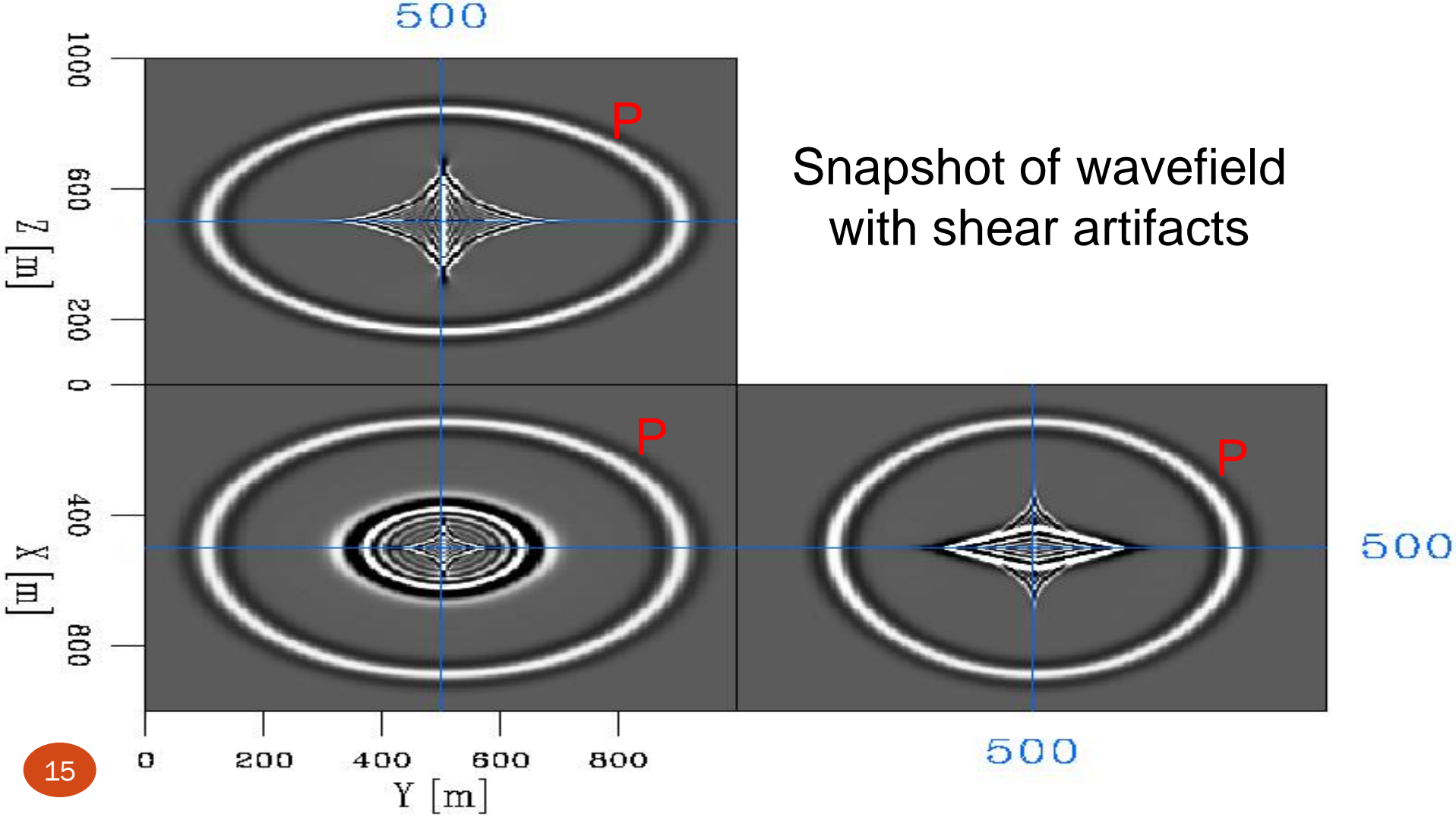
Wavenumber domain method

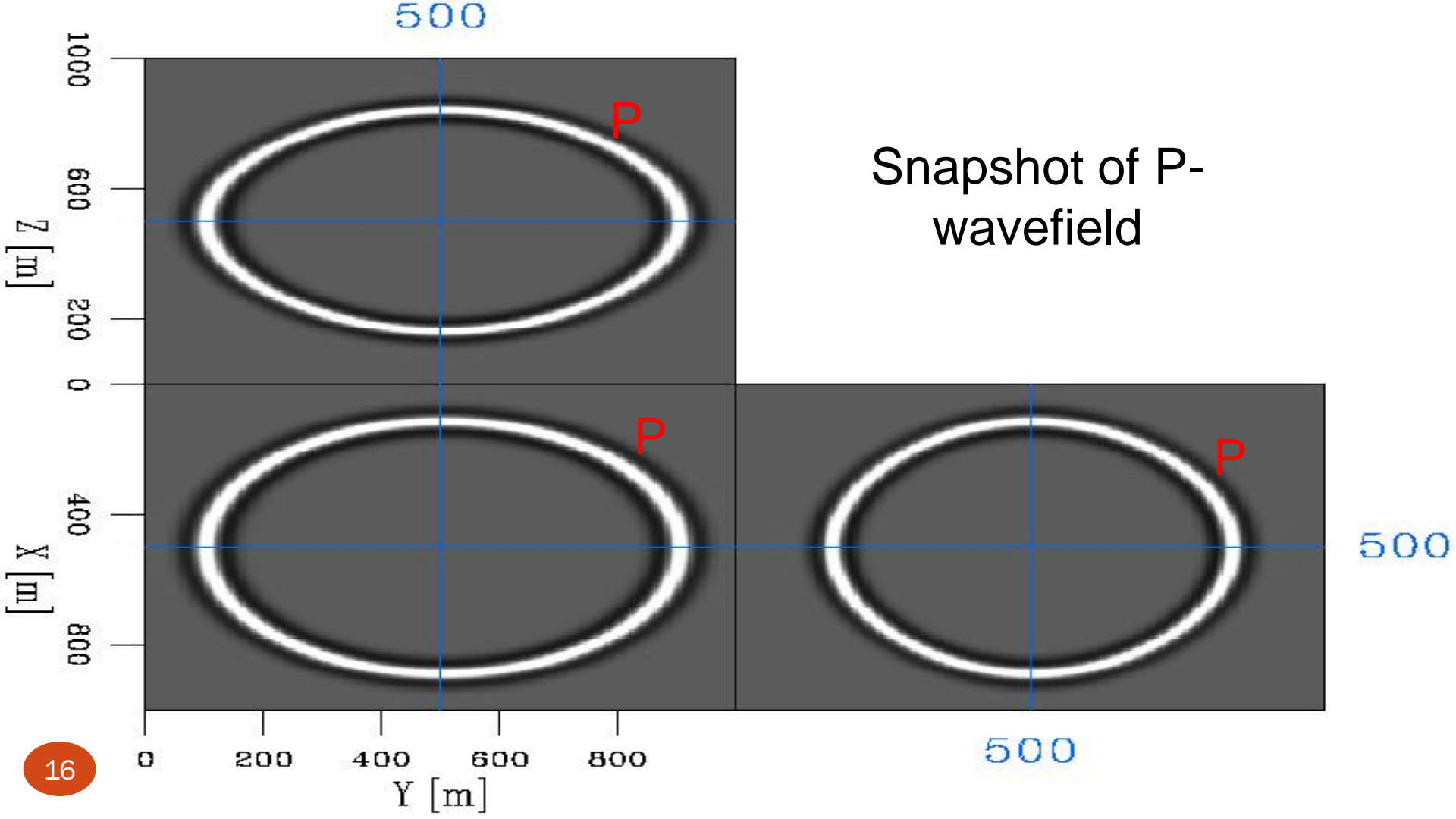
For all k 's: $M\tilde{\mathbf{D}} = Q\tilde{\Lambda}Q^{-1}, \tilde{\Lambda} = \text{diag} [\tilde{\lambda}_i(k)]$

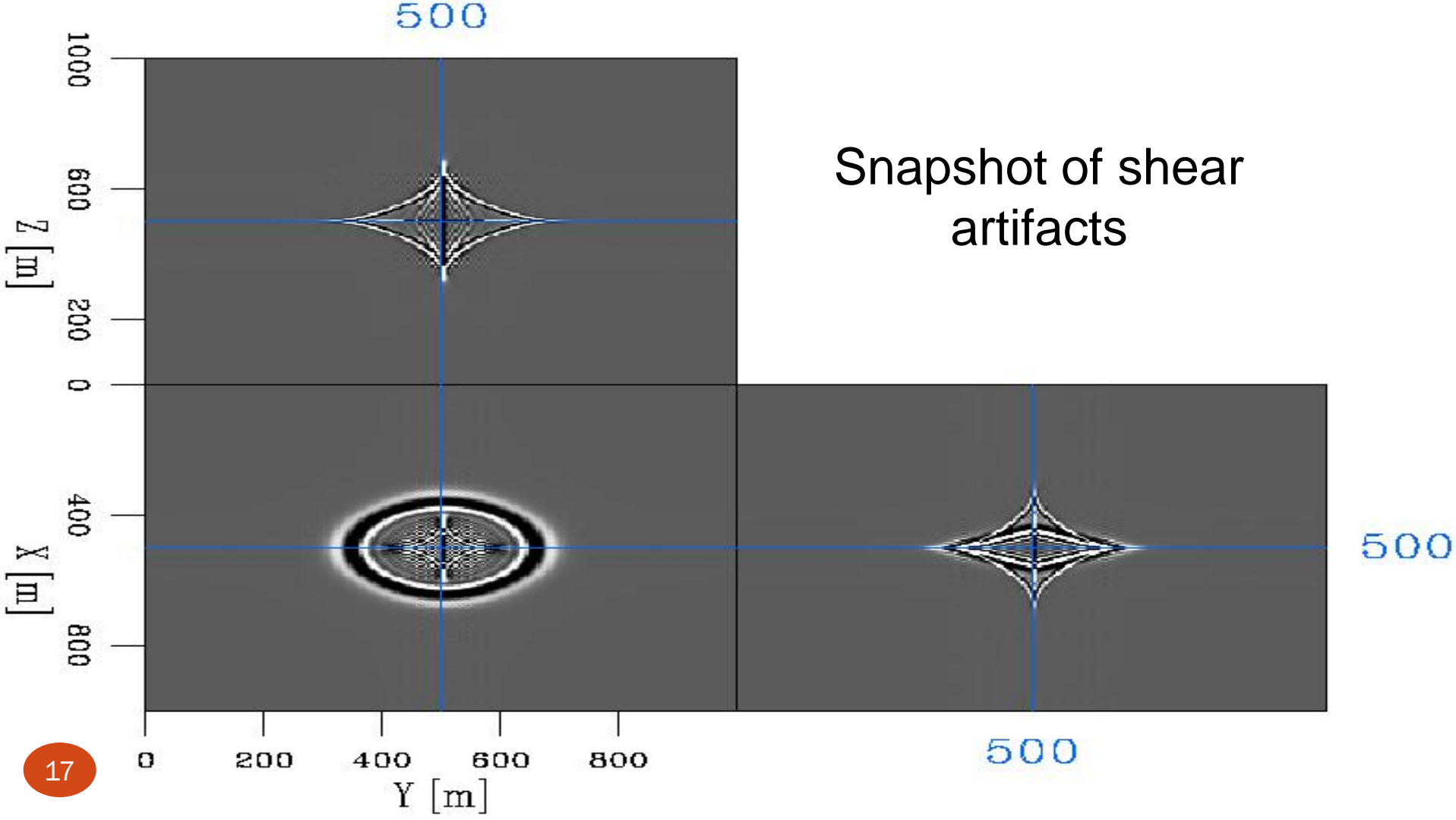
Choose eigenvalue corresponding to P-wave

Option 1: $p(x) \rightarrow \tilde{p}(k)$

$$\frac{\partial^2 \tilde{p}}{\partial t^2} = \tilde{\lambda}_i(k) \tilde{p}(k)$$







Space-time domain method

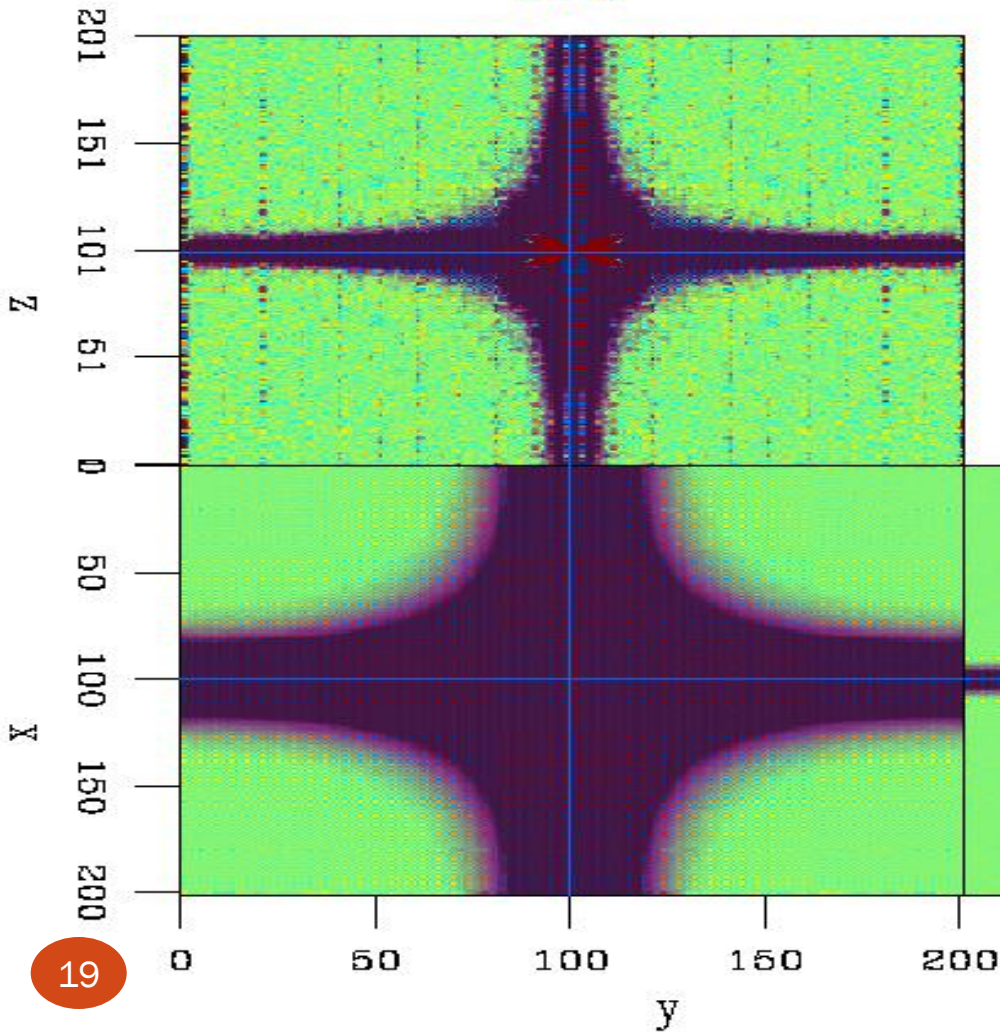
For all k 's: $M\tilde{\mathbf{D}} = Q\tilde{\Lambda}Q^{-1}, \tilde{\Lambda} = \text{diag}[\tilde{\lambda}_i(k)]$

Choose eigenvalue corresponding to P-wave

Option 2: $\tilde{\lambda}_i(k) \rightarrow \lambda_i(x)$

$$\frac{\partial^2 p}{\partial t^2} = \lambda_i(x) * p(x)$$

100

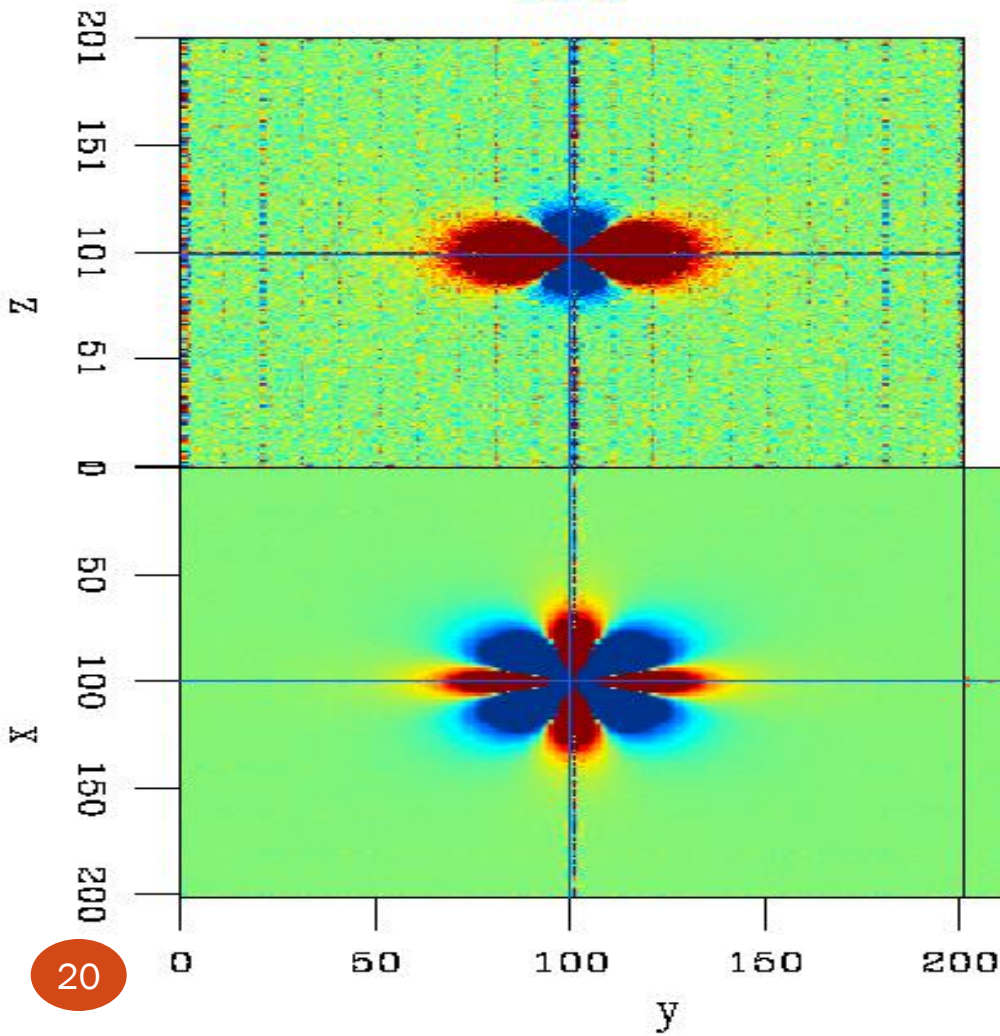


Factorizing exact
Fourier representation

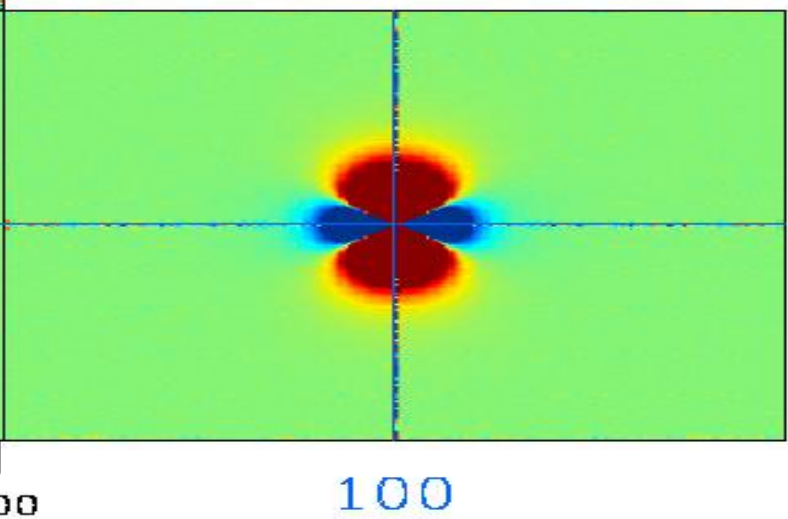
100

100

100

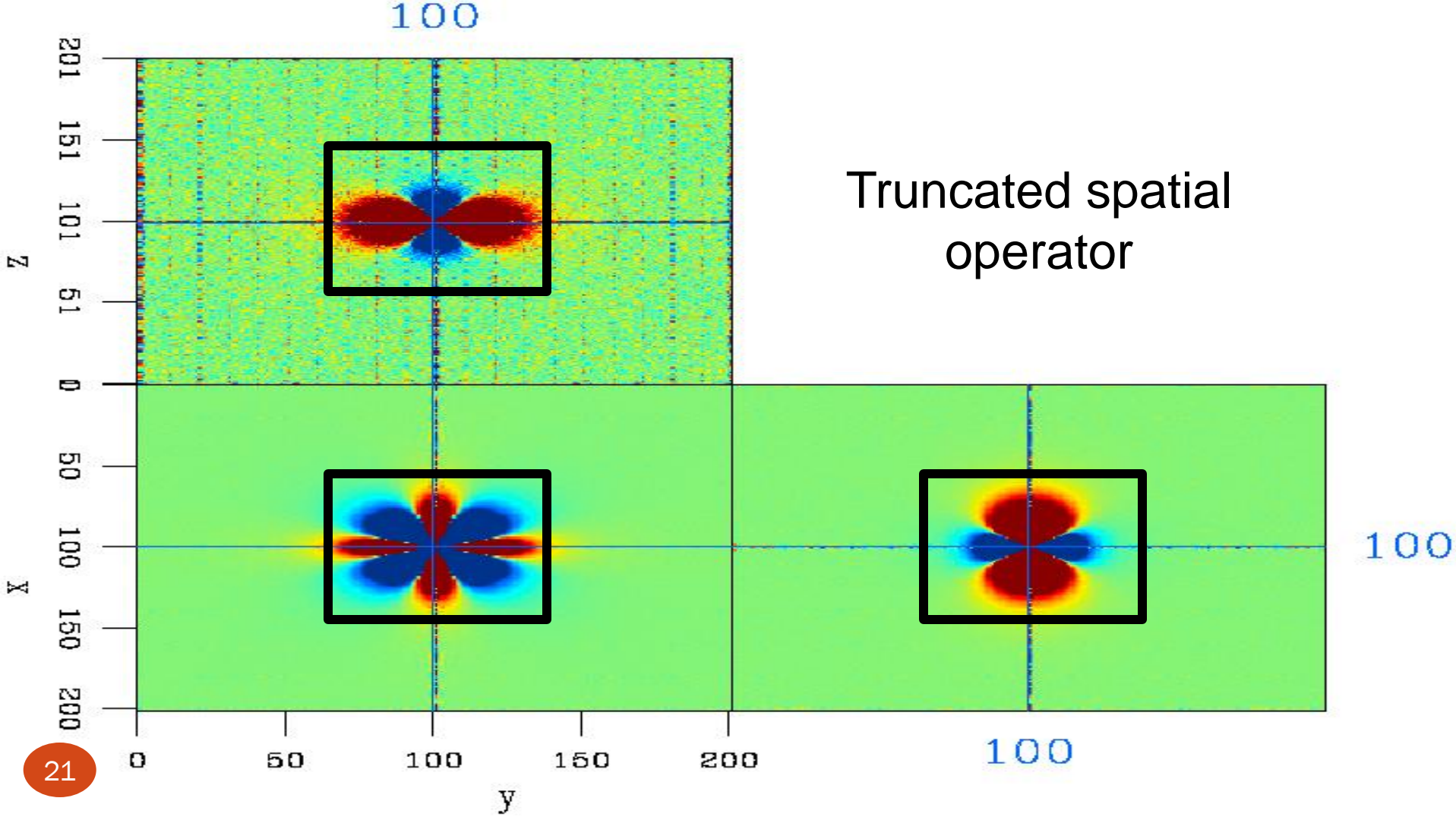


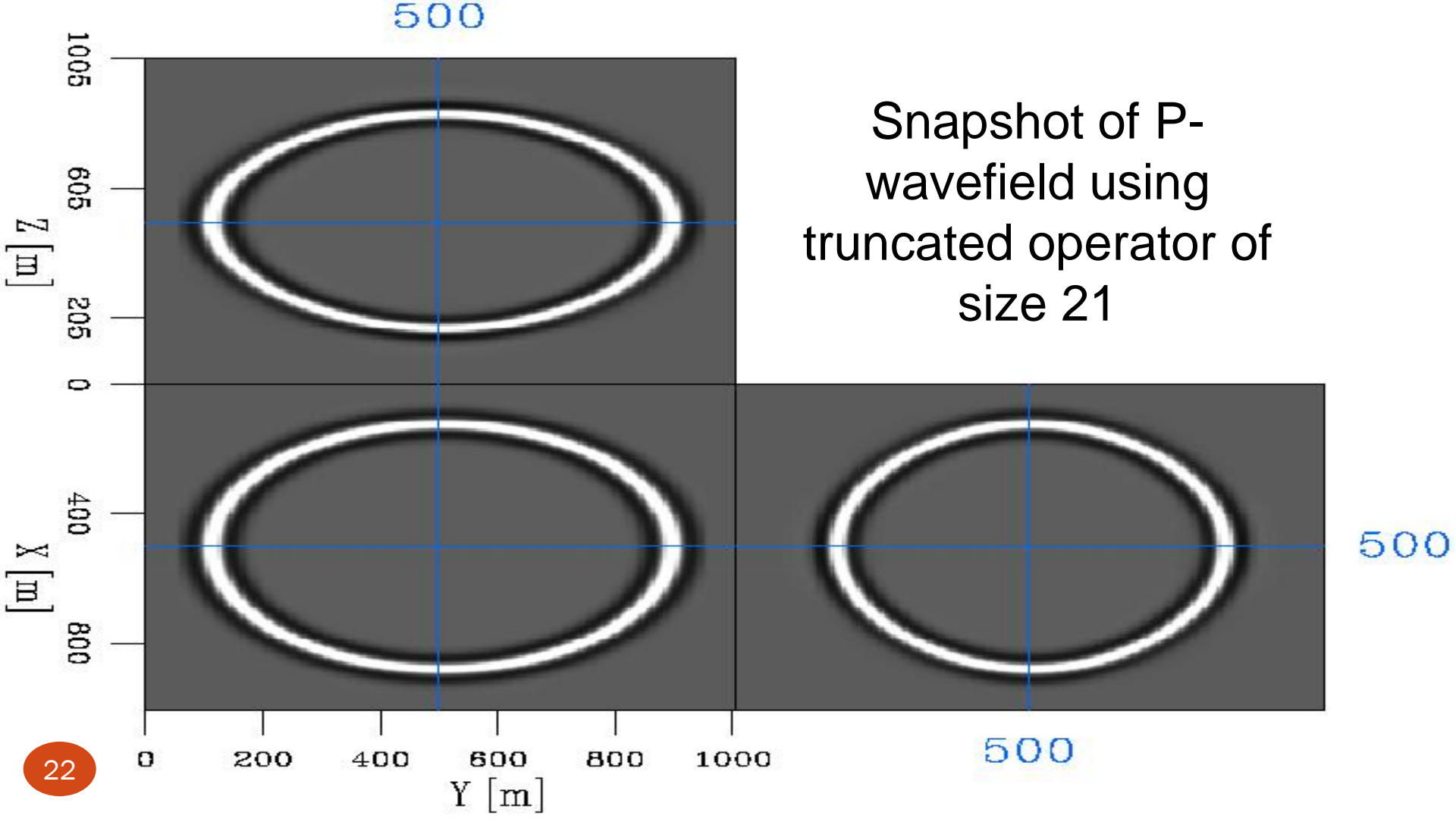
Factorizing FD
approximation

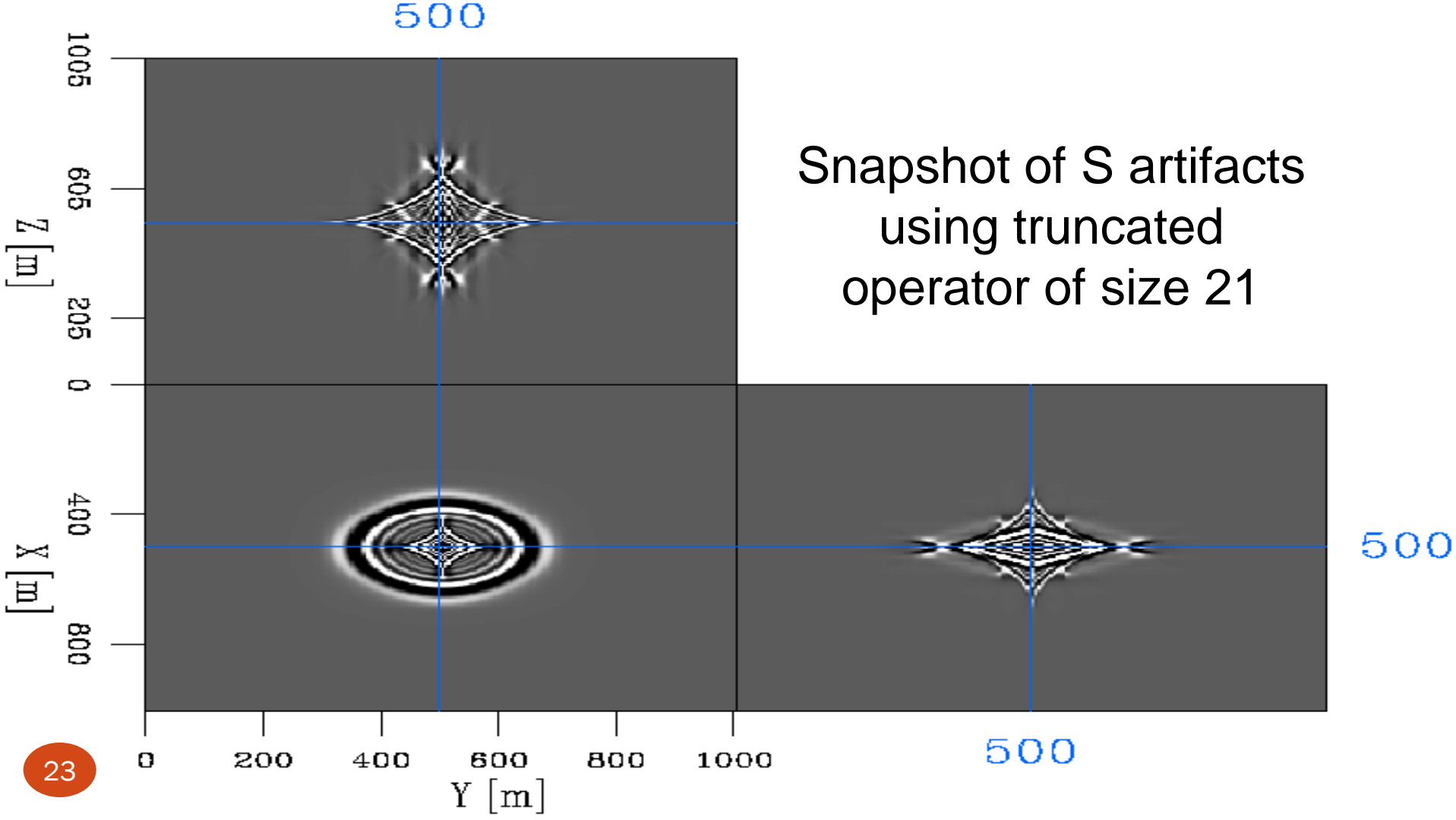


100

100







Operator's accuracy after truncation

- Wavenumber domain: accurate but slow
- Spatial domain: less accurate but fast
- Both use 2nd-order temporal FD, same time steps, and spatial discretizations
 - => same temporal dispersion error
- Spatial subtracts wavenumber
 - => spatial dispersion error

Spatial dispersion error

Isotropic, star stencil

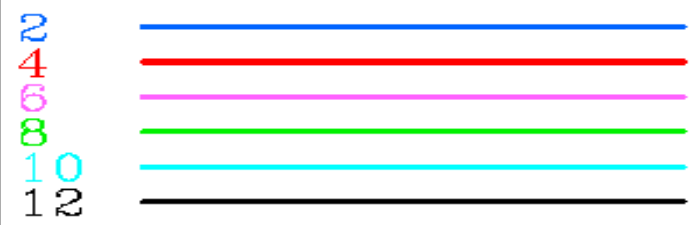
amplitude

0.1

0

-0.1

Stencil size



0

0.04

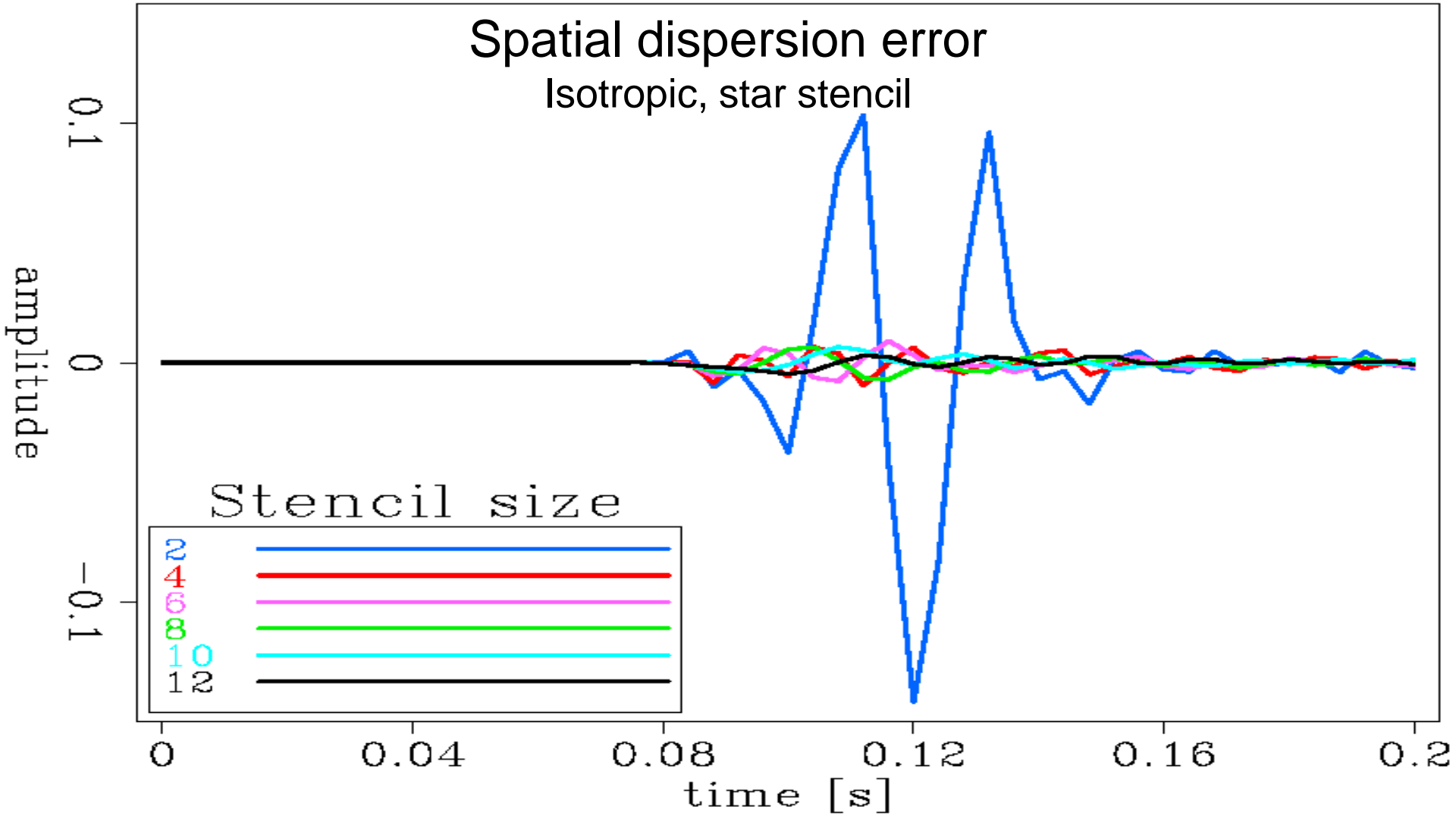
0.08

0.12

0.16

0.2

time [s]



Spatial dispersion error

Isotropic, star stencil

amplitude

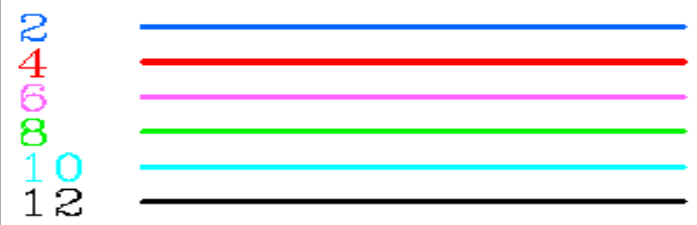
0.1

0

-0.1

Acceptable level of accuracy

Stencil size



0

0.04

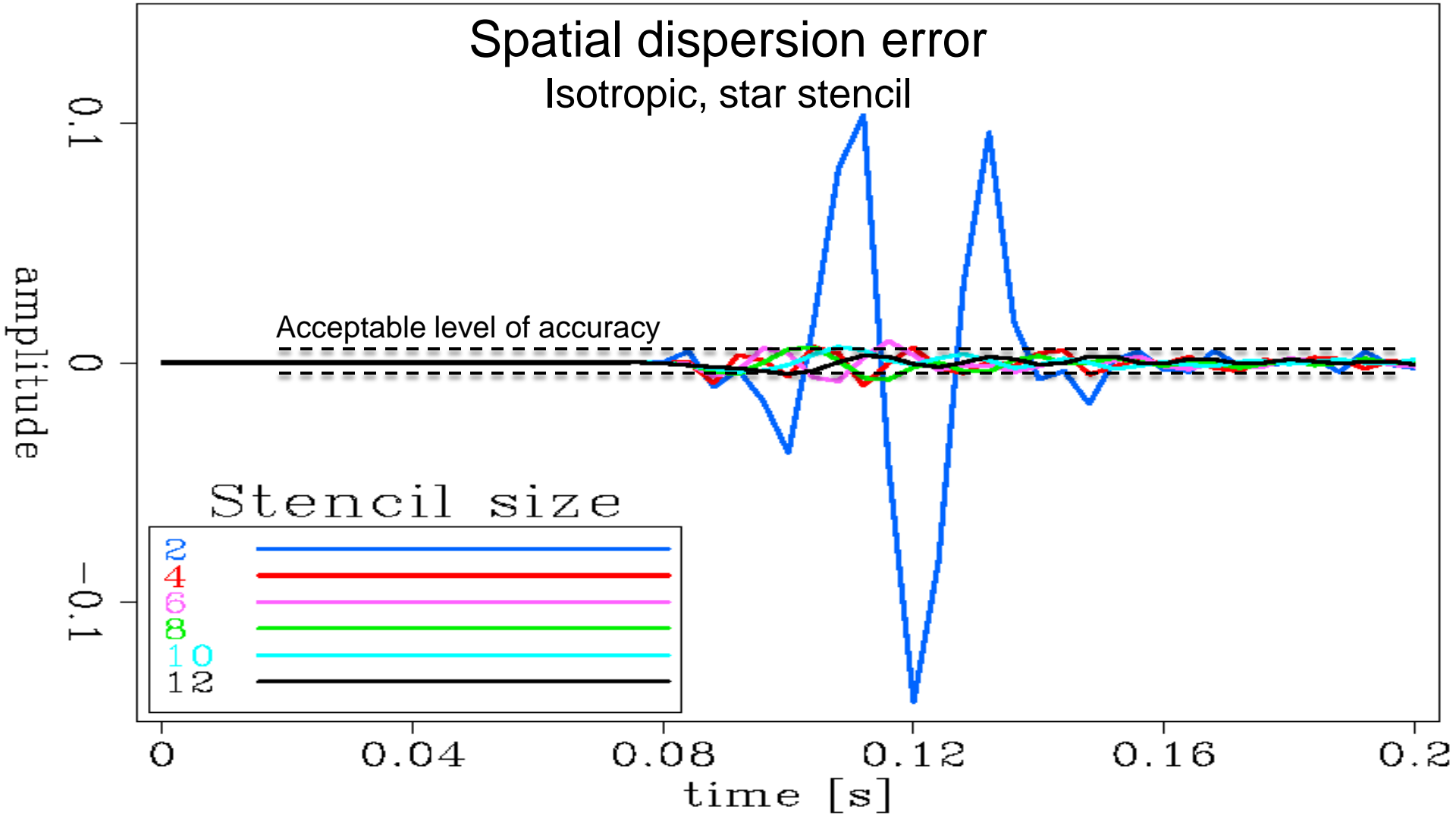
0.08

0.12

0.16

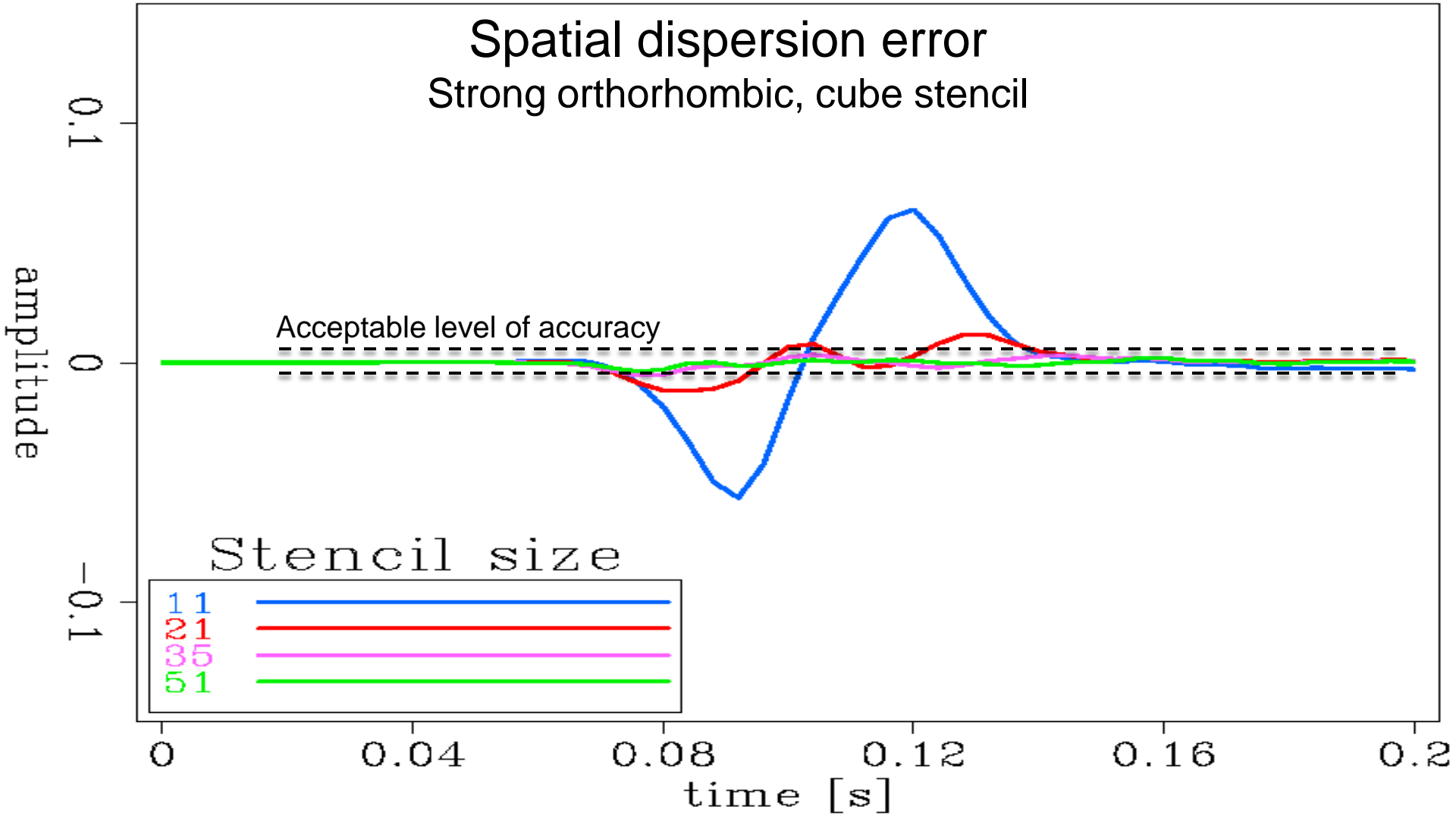
0.2

time [s]



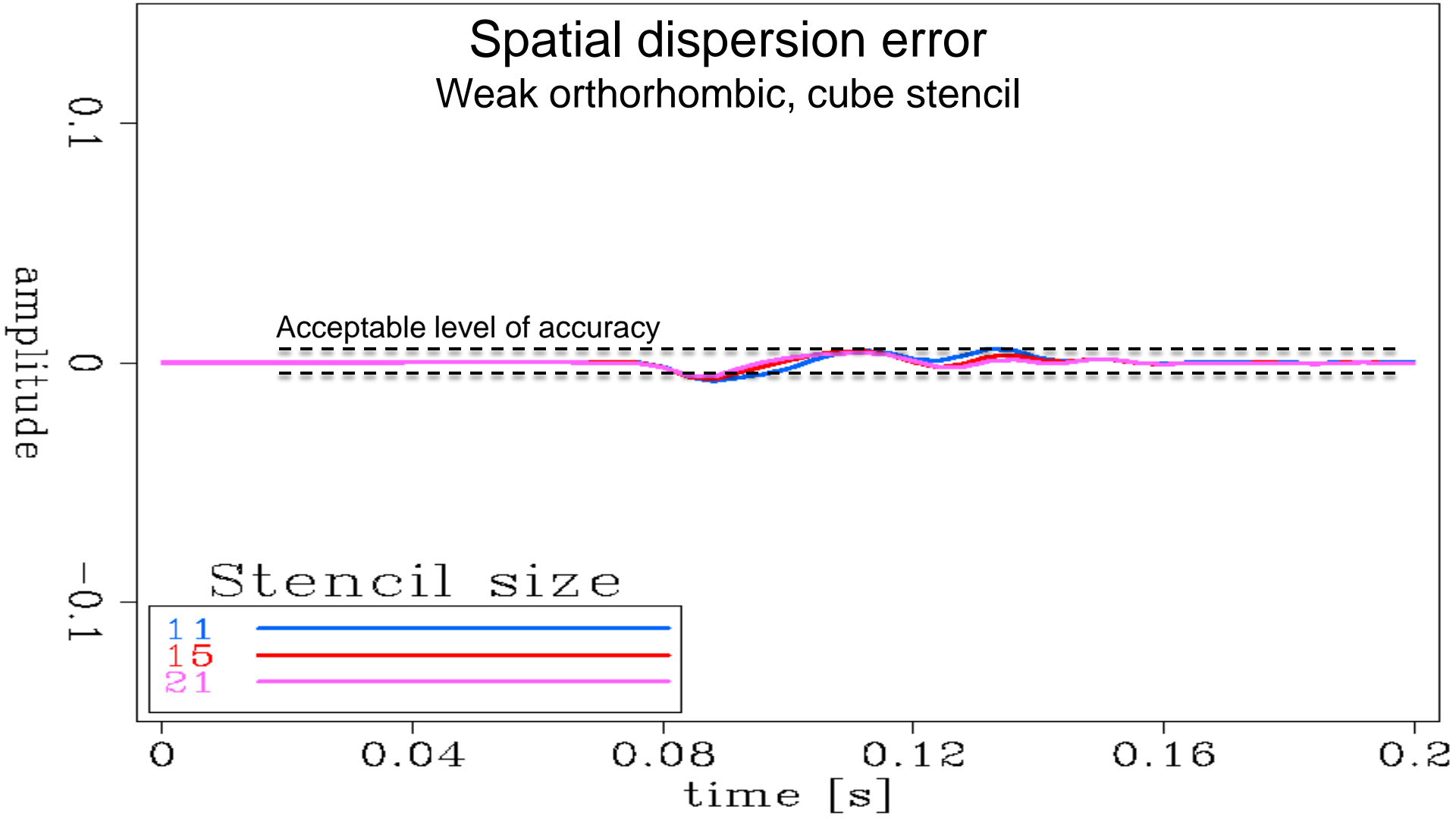
Spatial dispersion error

Strong orthorhombic, cube stencil



Spatial dispersion error

Weak orthorhombic, cube stencil



amplitude

Acceptable level of accuracy

Stencil size

11
15
21

time [s]

Spatial dispersion error

VTI, cube stencil

amplitude

0.1

0

-0.1

Acceptable level of accuracy

Stencil size

11
21
35



0

0.04

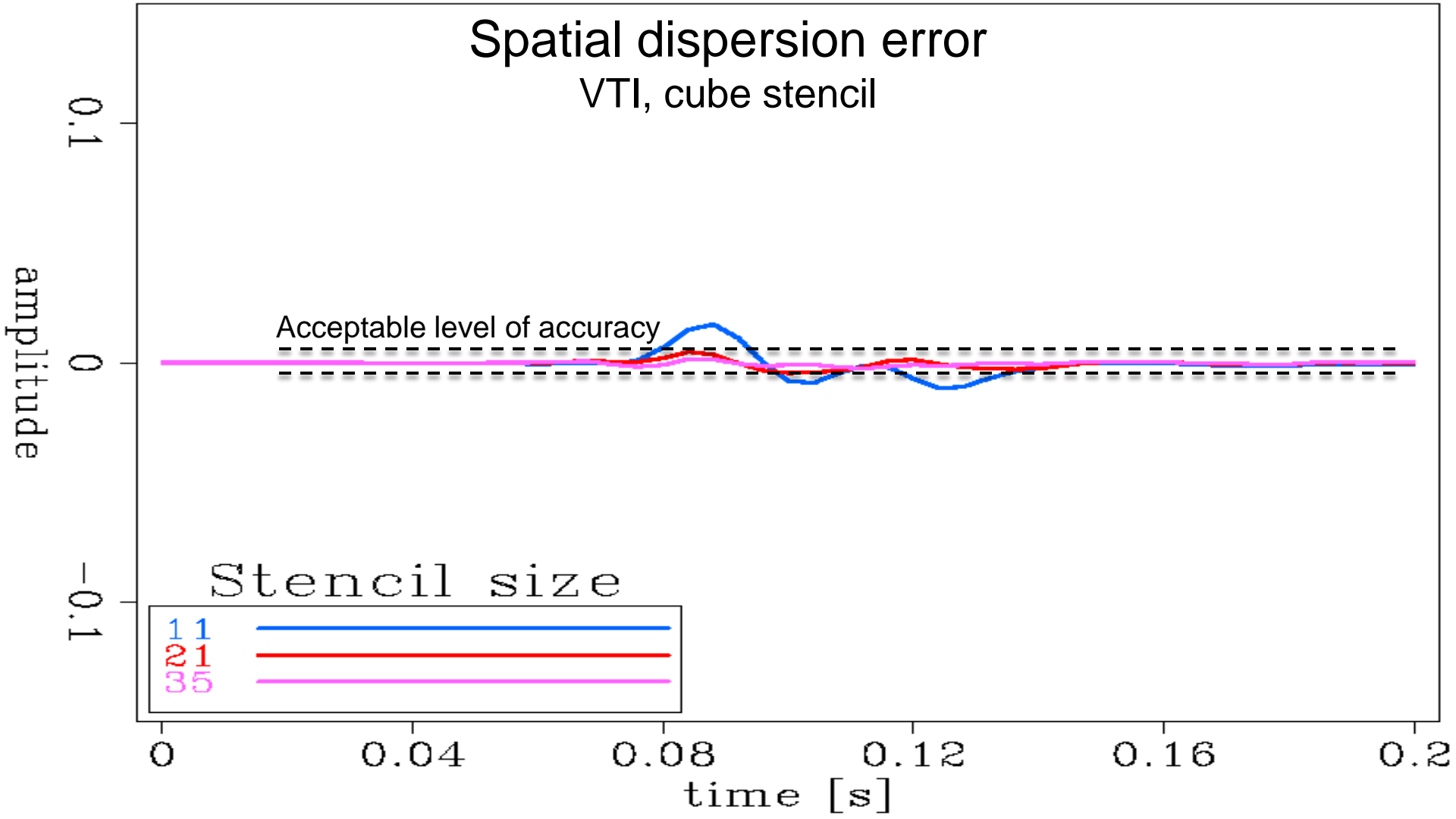
0.08

0.12

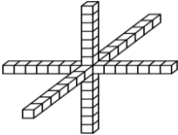
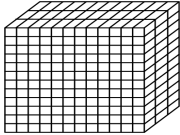
0.16

0.2

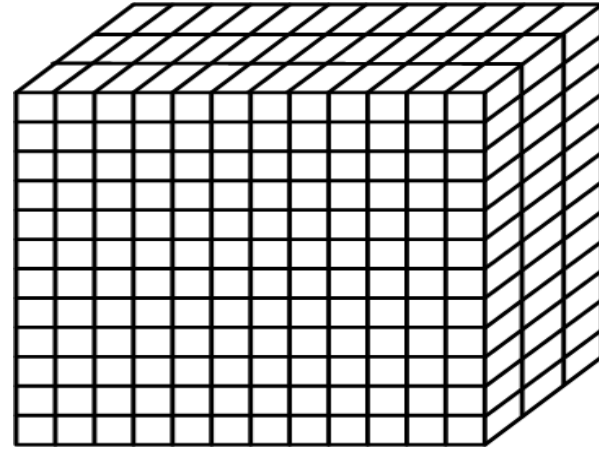
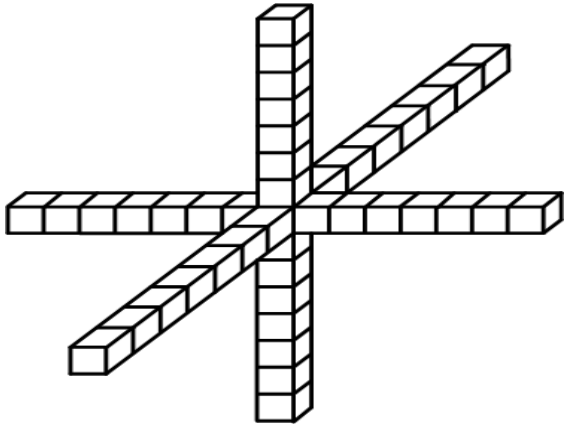
time [s]



Cost comparison

	$\frac{\partial^2}{\partial t^2} \boldsymbol{\sigma} = \mathbf{M} \mathbf{D} \boldsymbol{\sigma}$	$\frac{\partial^2 p}{\partial t^2} = \lambda_i(x) * p(x)$
Stored wavefields	3	1
Computation (+ and \times)	1 (10 th -order FD) 	140 (21-sized cube) 

Cost comparison

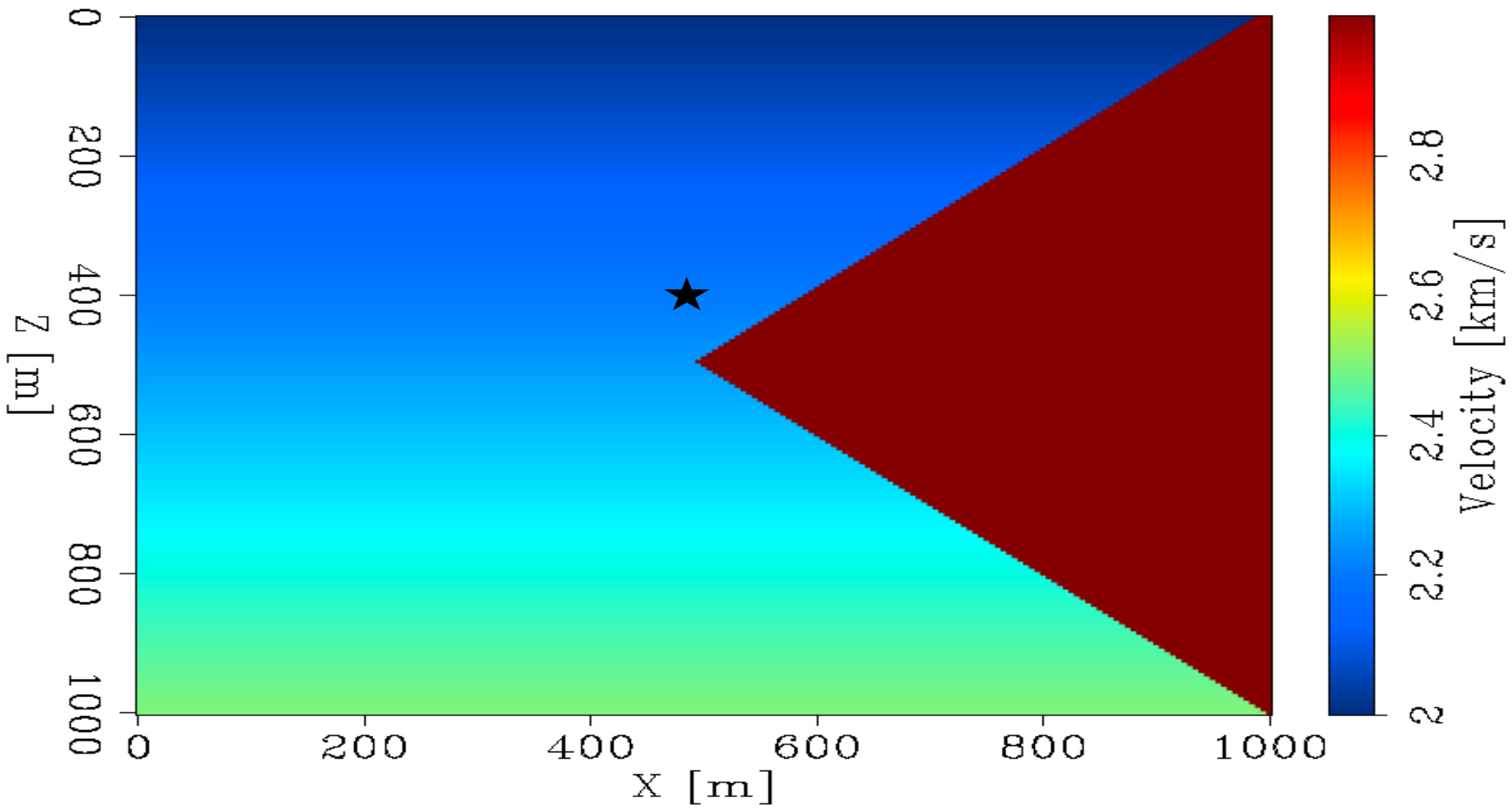


Computation hides latency
Highly parallelizable
Better cache reuse

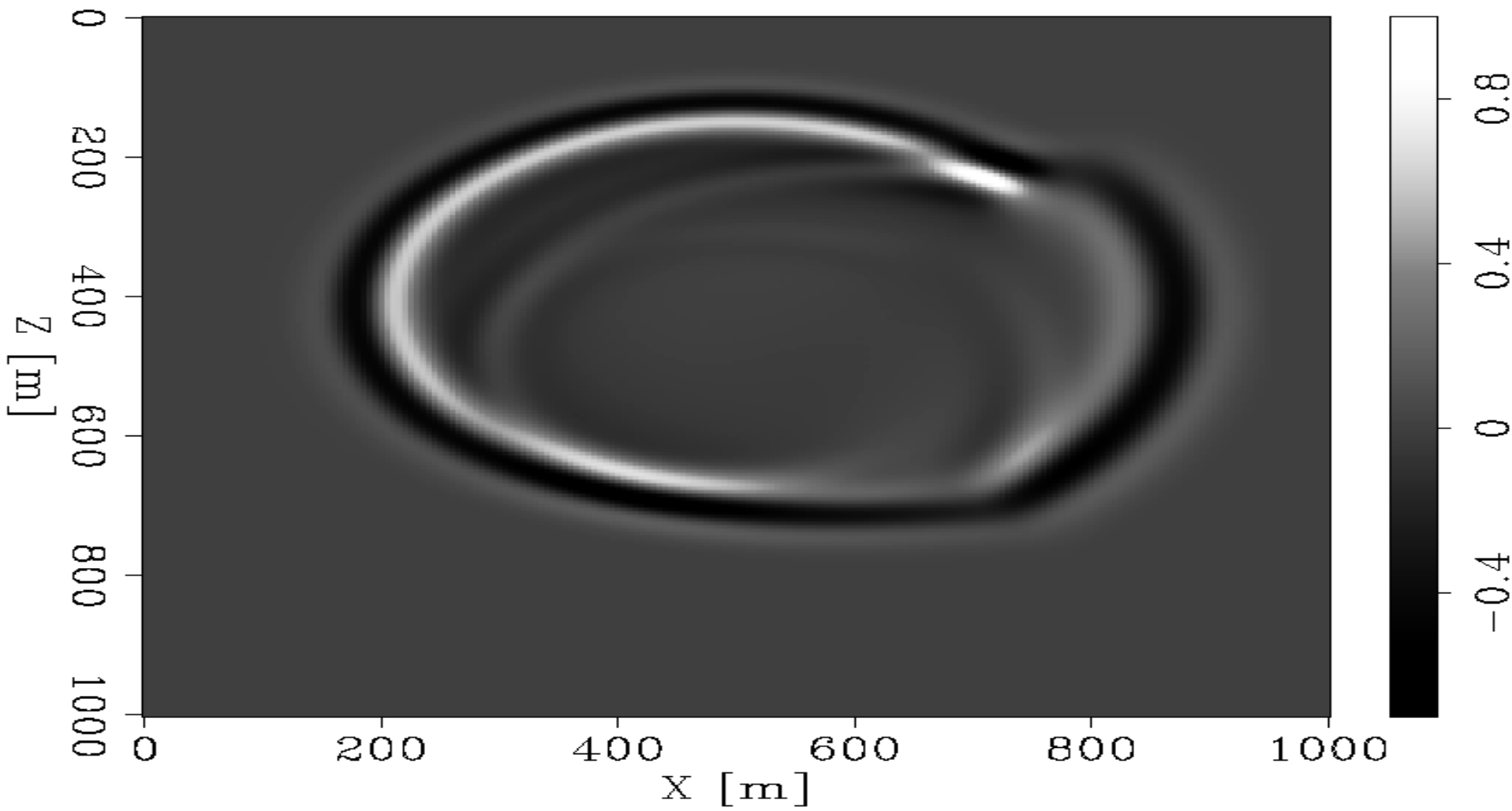
Application to inhomogeneous media

- Spatially varying operator
- Exact solution requires calculation of operator for **every** present model parameters.

True model



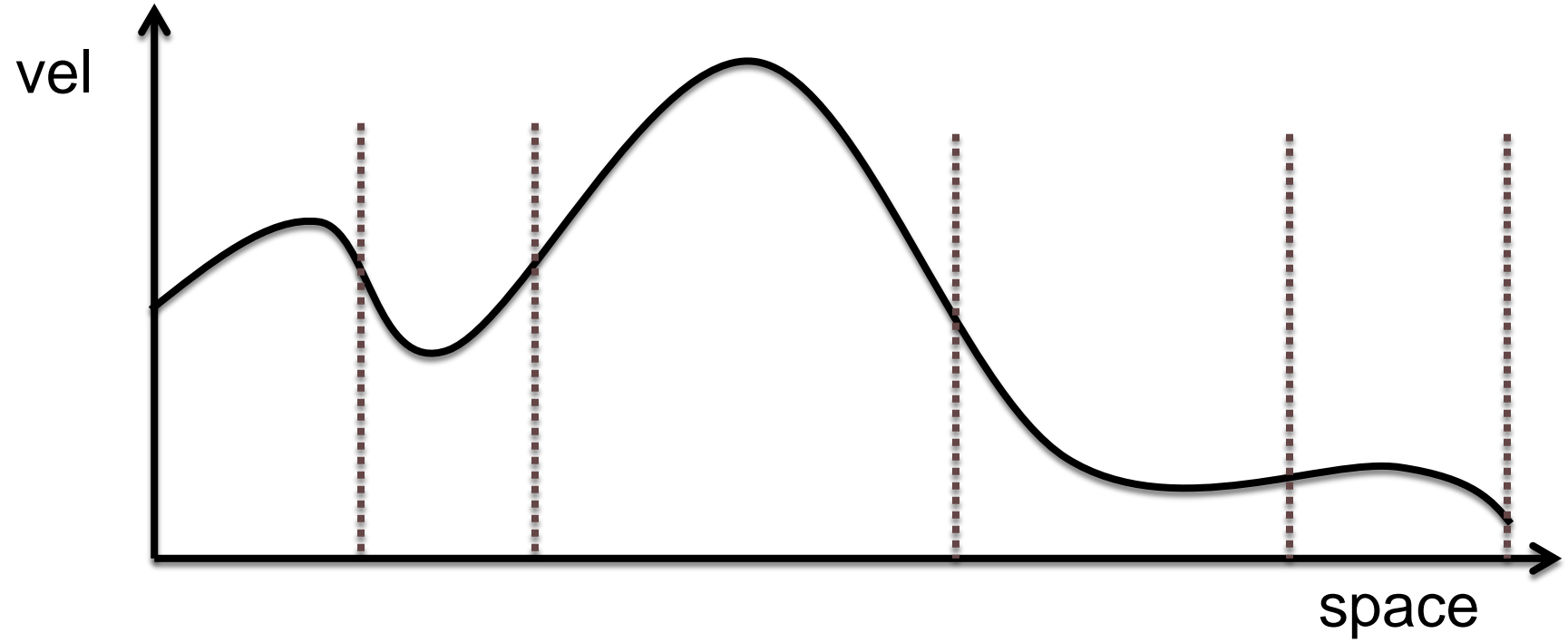
True wavefield



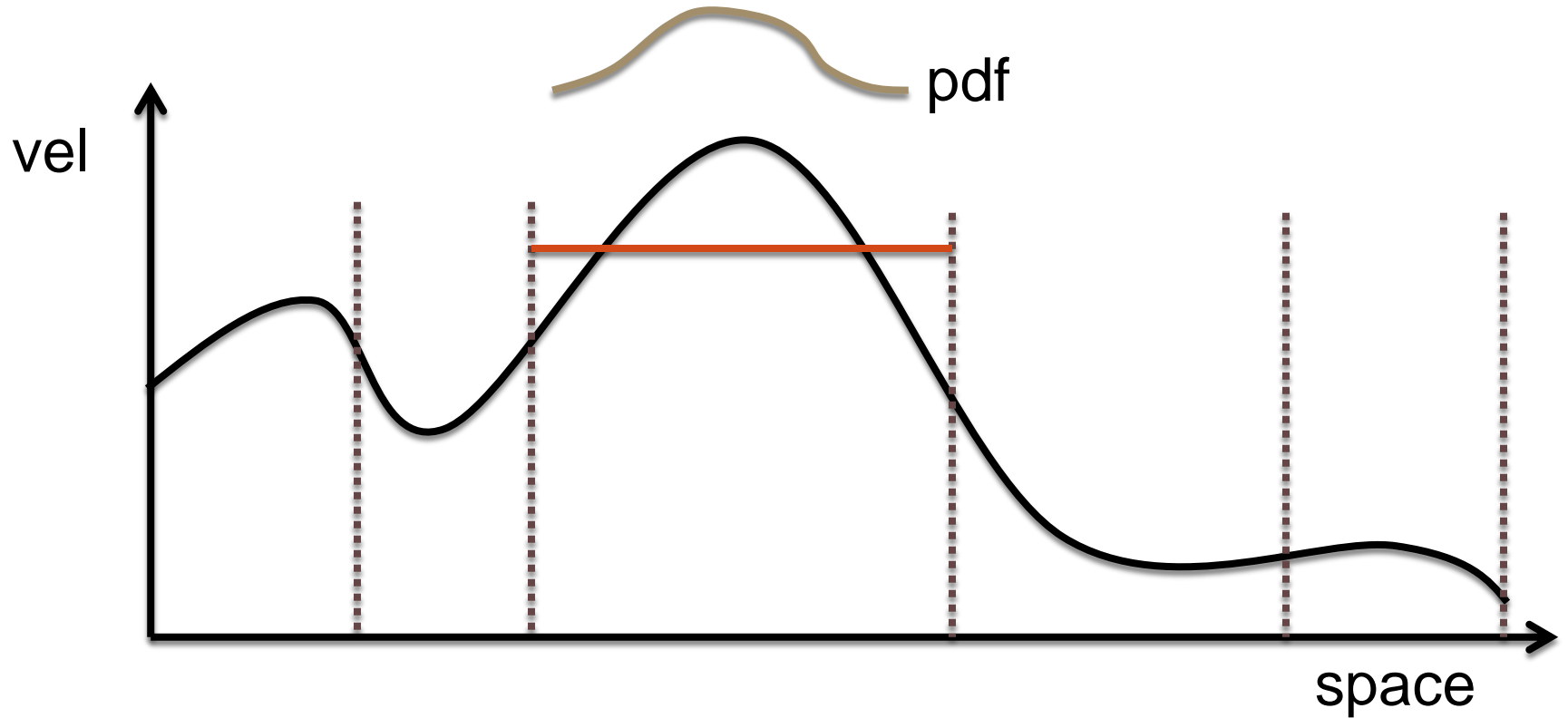
Application to inhomogeneous media

- Spatially varying operator
- Exact solution requires calculation of operator for **every** present model parameters
 - => expensive
- Approximation: compute operators for a number of reference model parameters
 - Reference selection: Lloyd algorithm (Clapp, 2006)
- Interpolation of operators

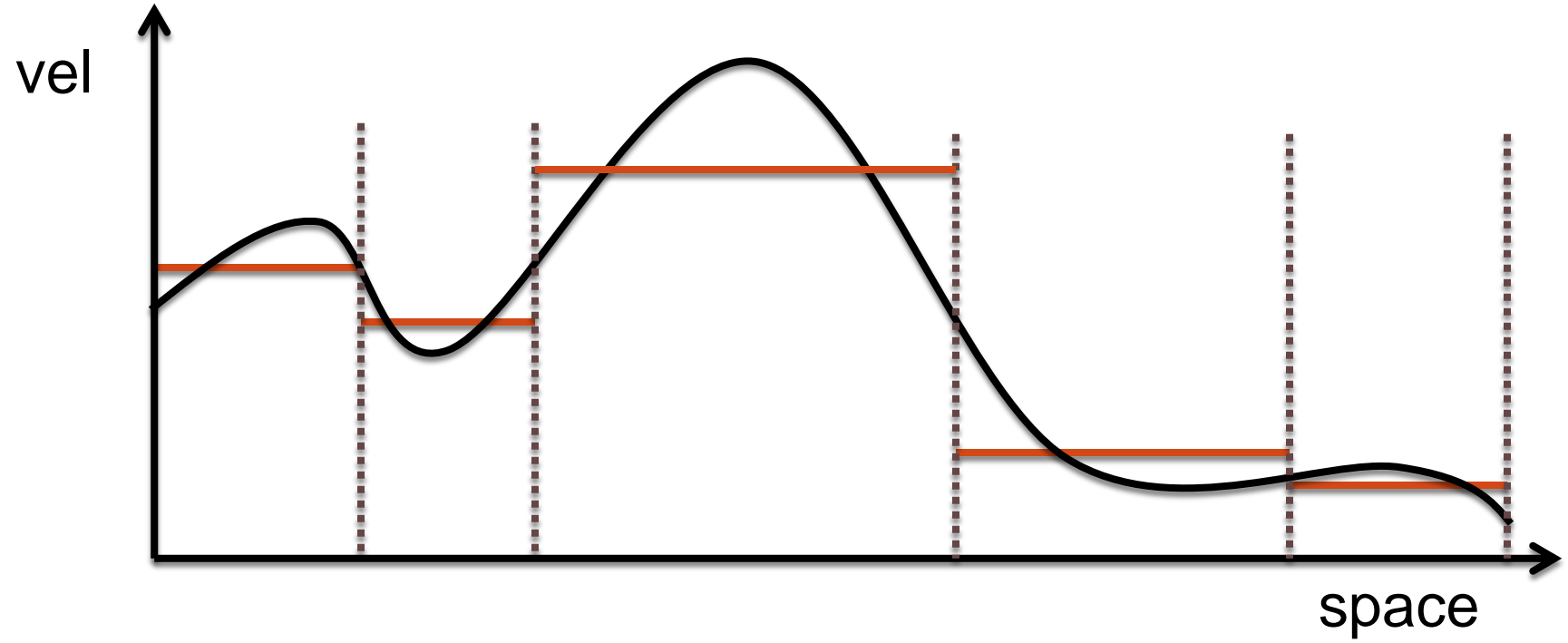
Lloyd algorithm for reference selection



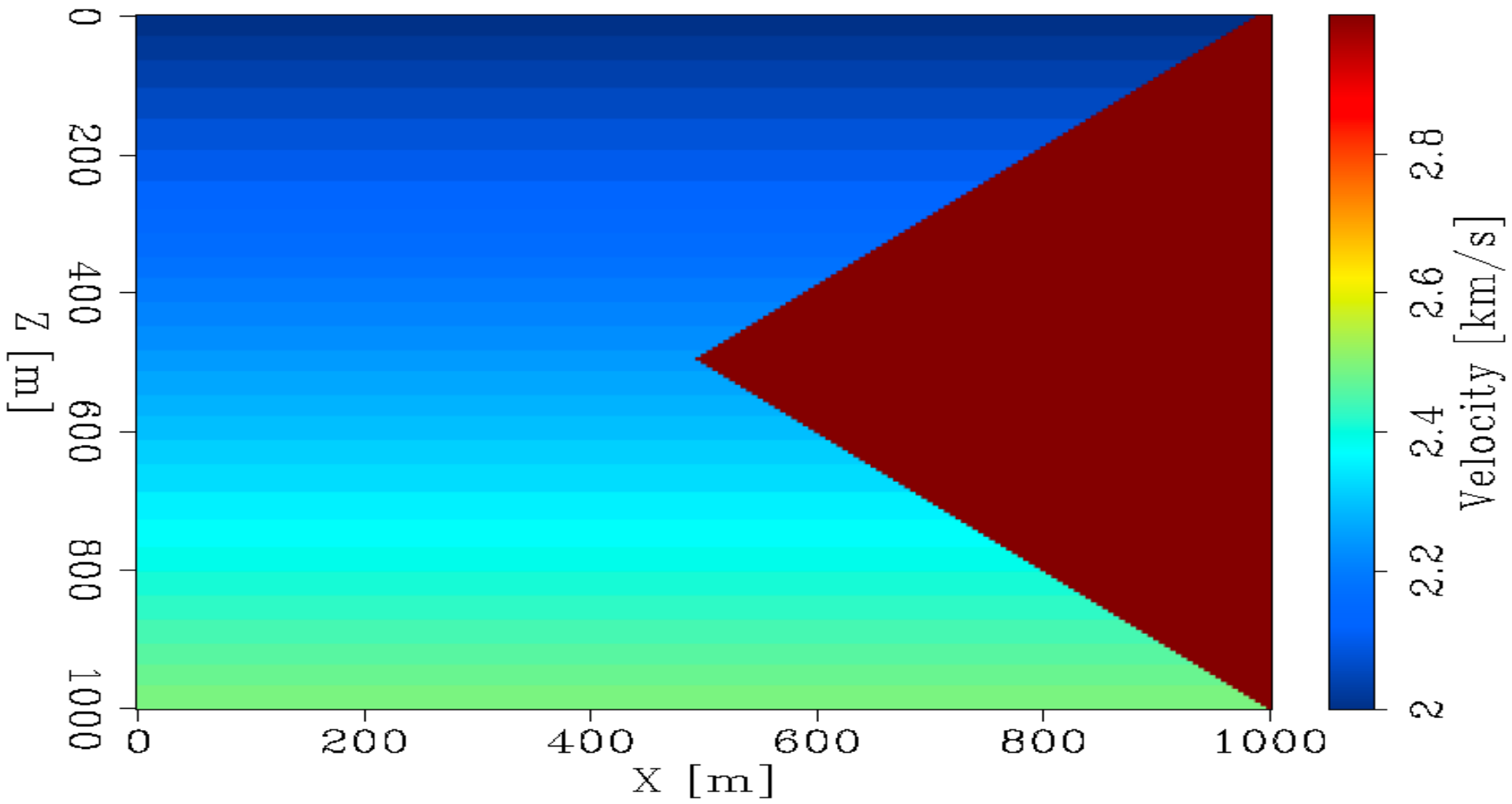
Lloyd algorithm for reference selection



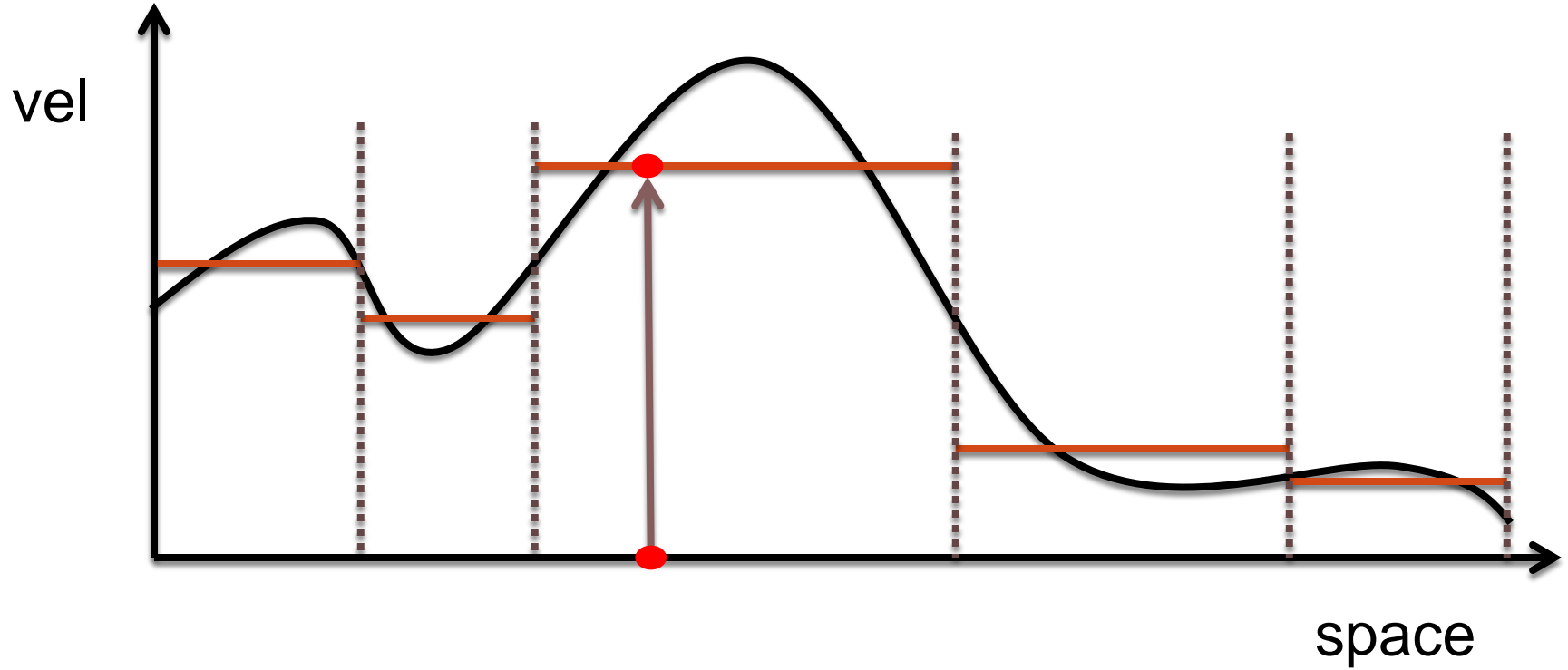
Lloyd algorithm for reference selection



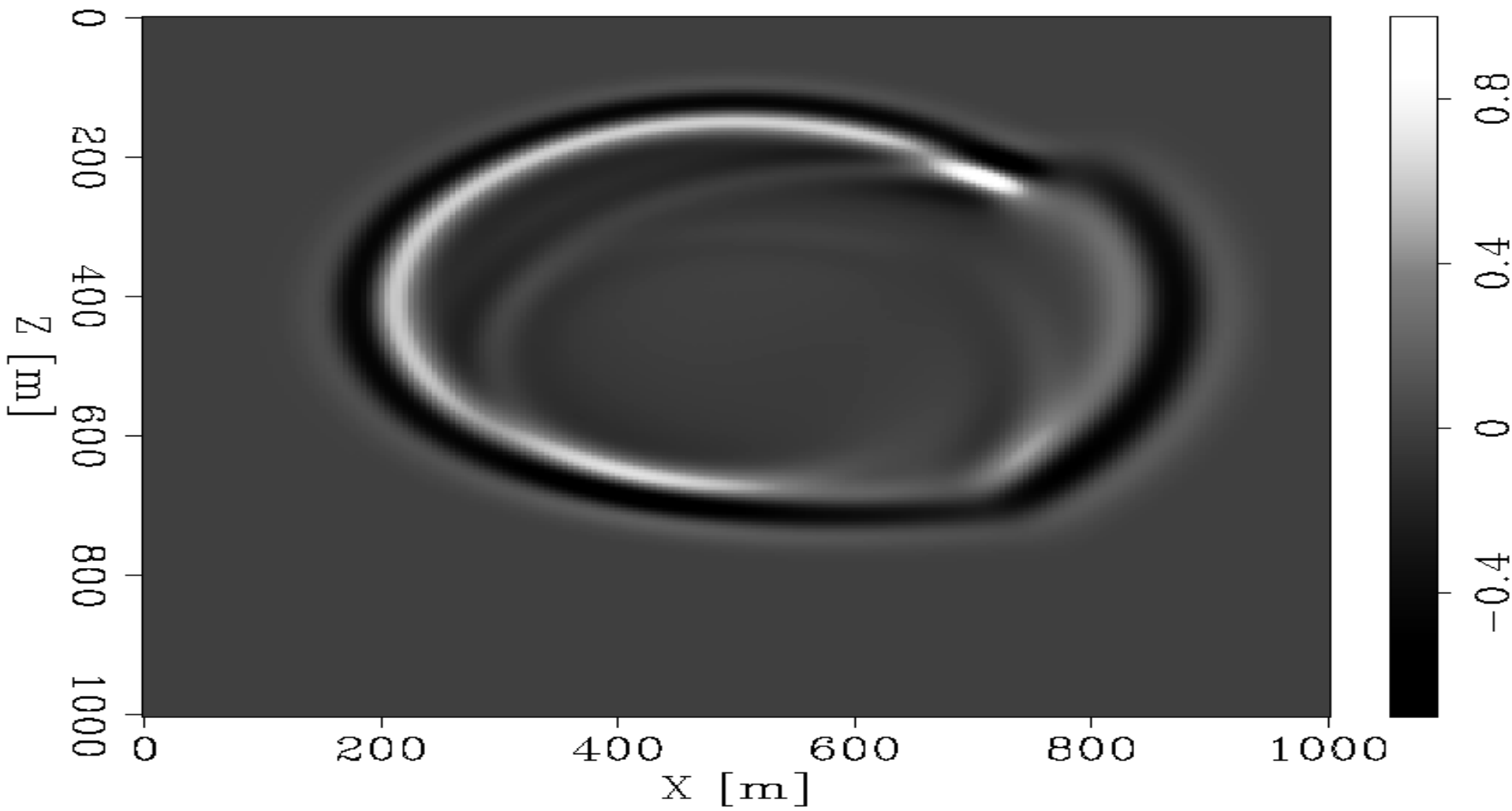
Model w/ 29 refs by Lloyd



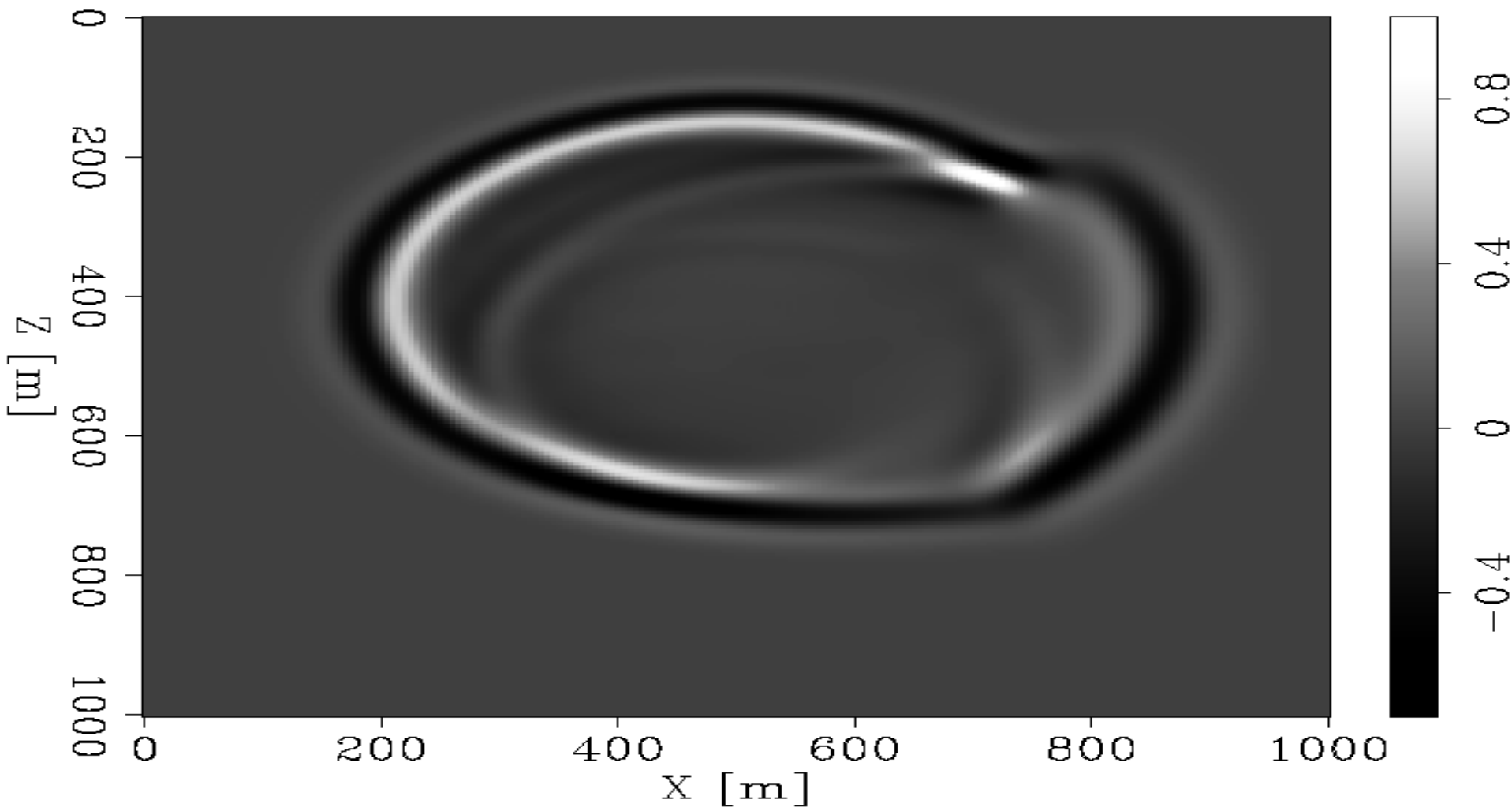
No interpolation



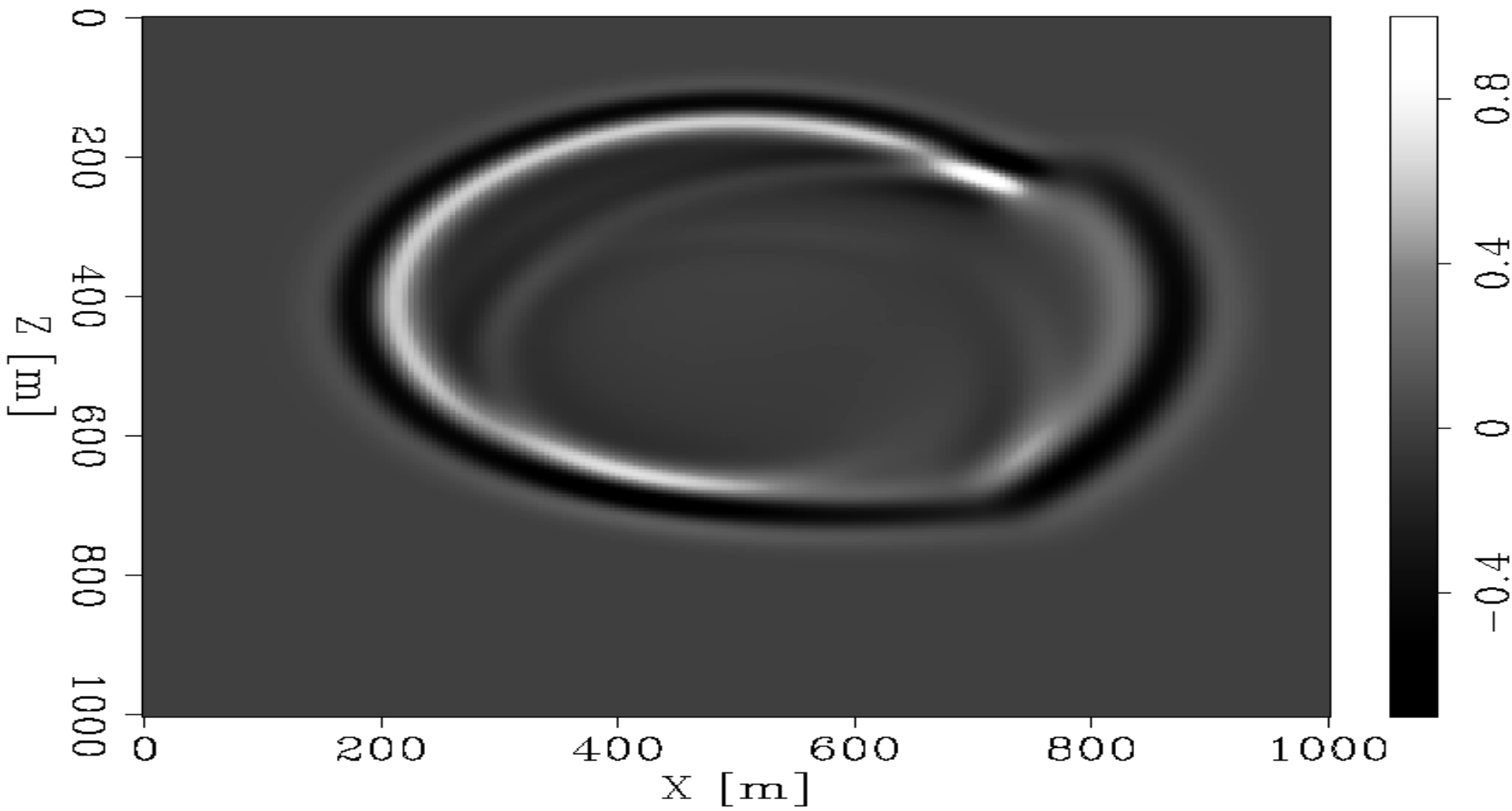
True wavefield



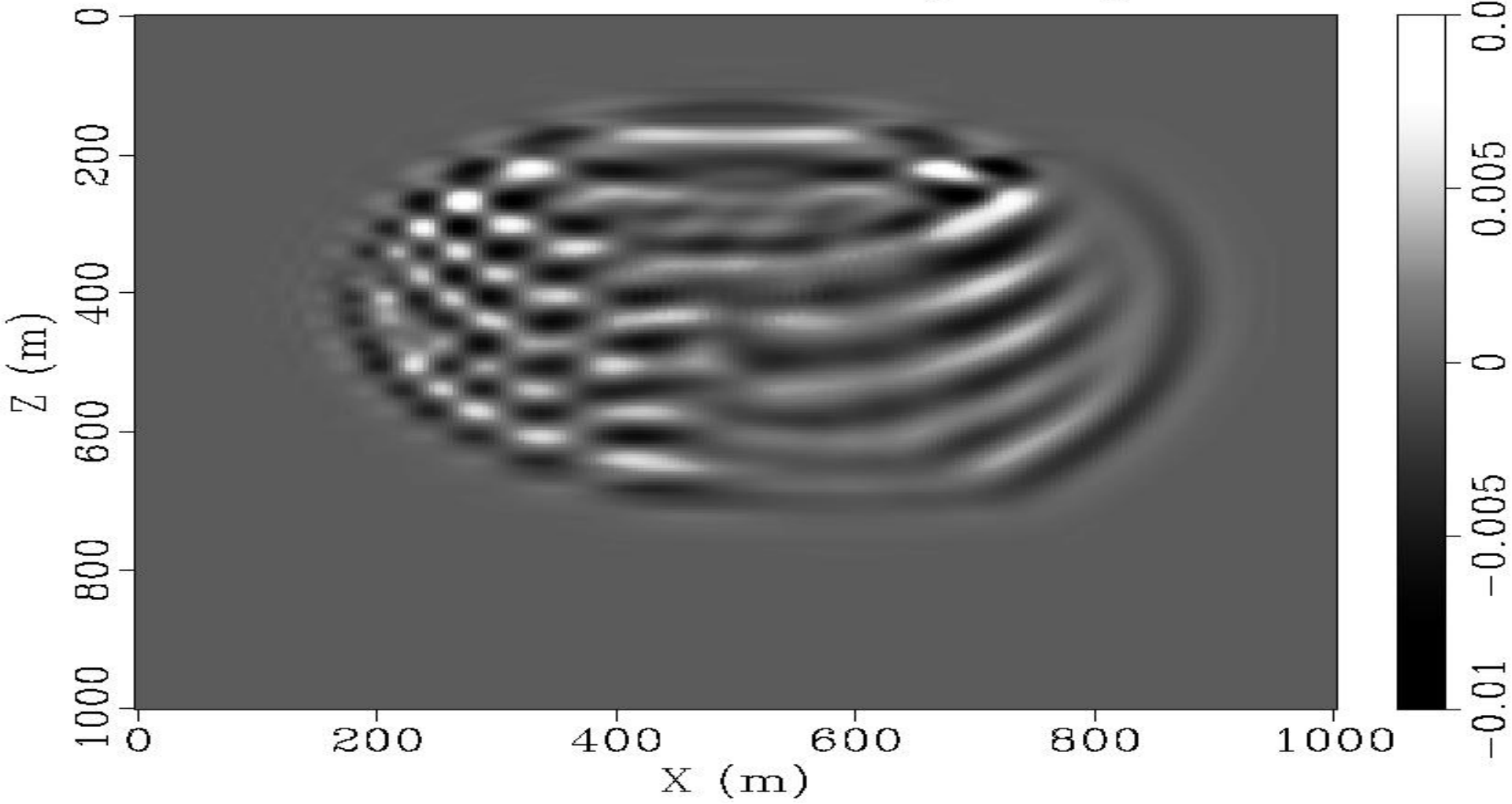
29 refs



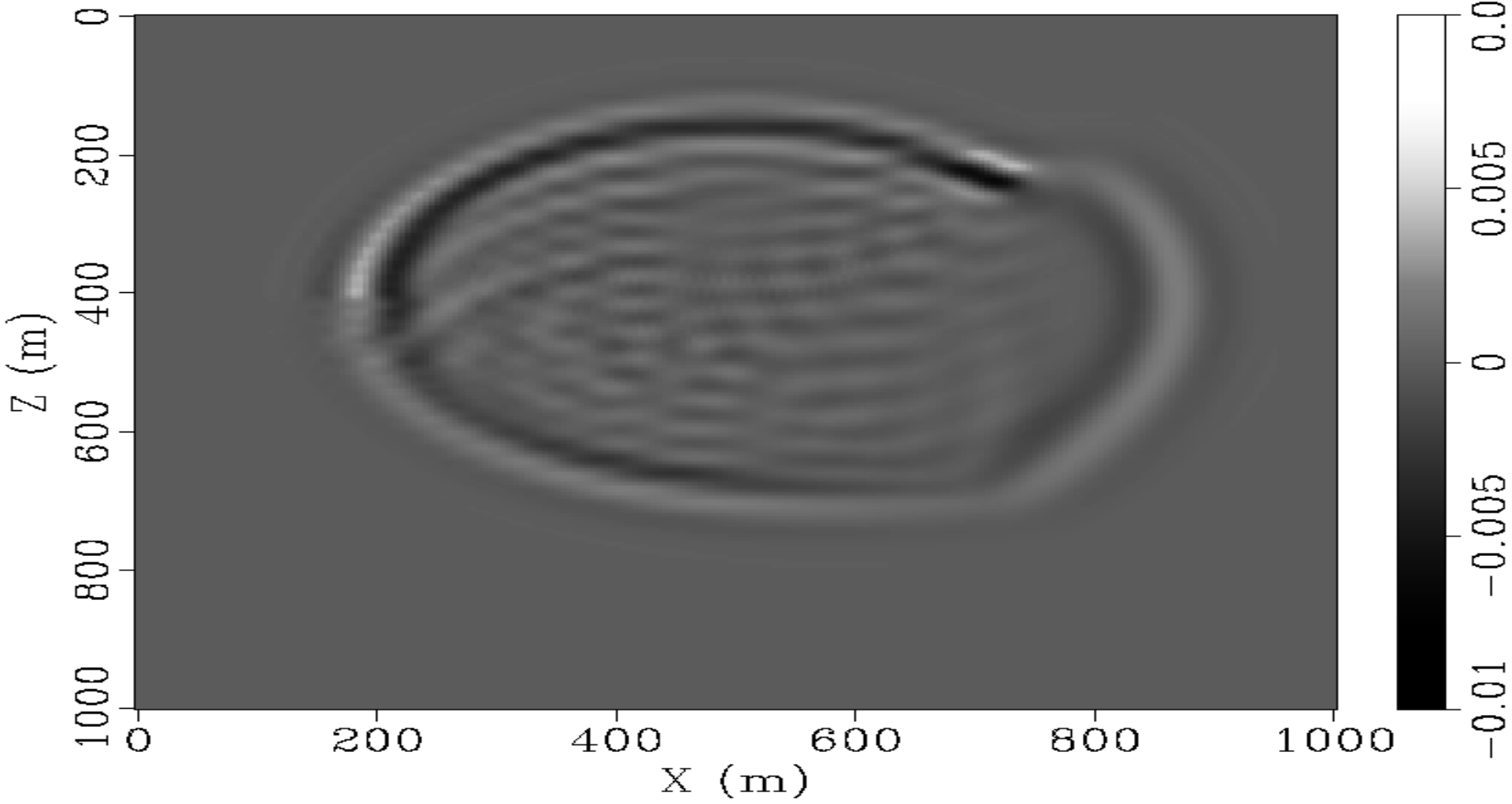
51 refs



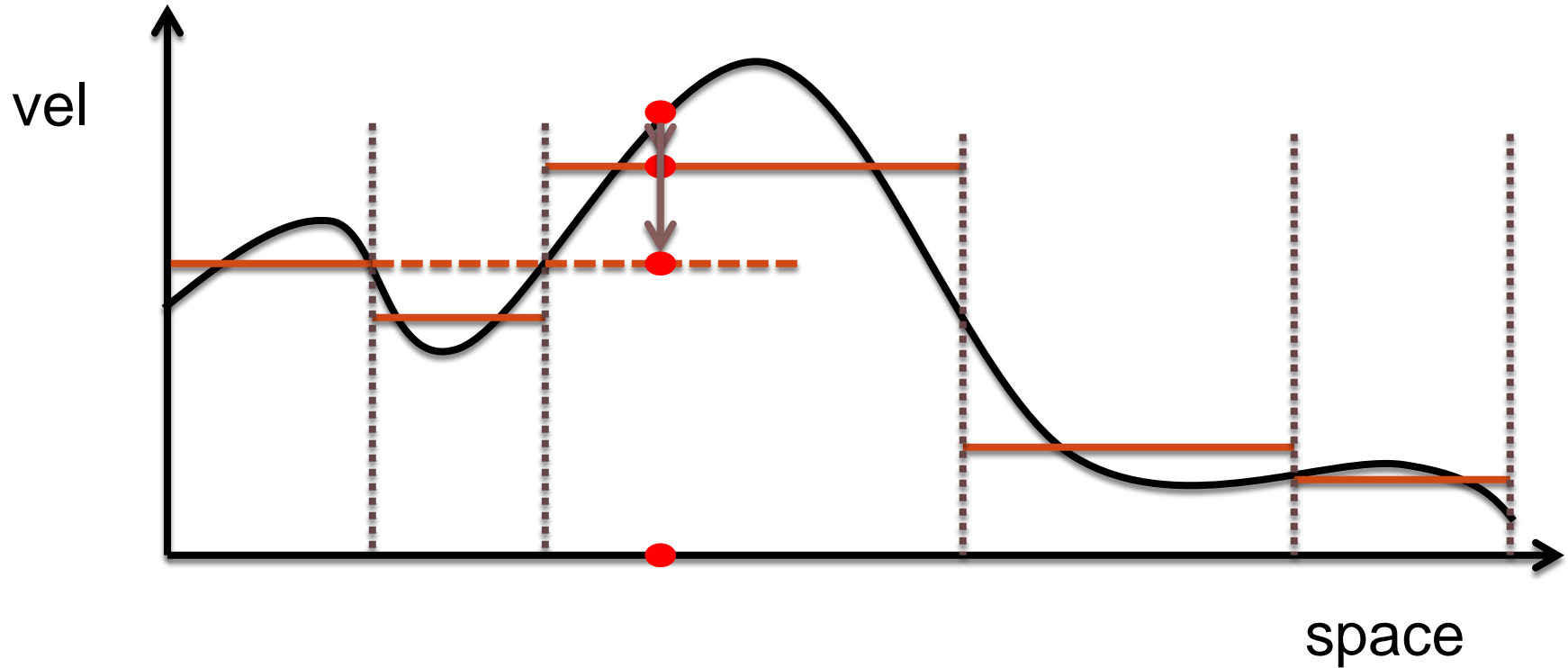
Error of 29 refs (100x)



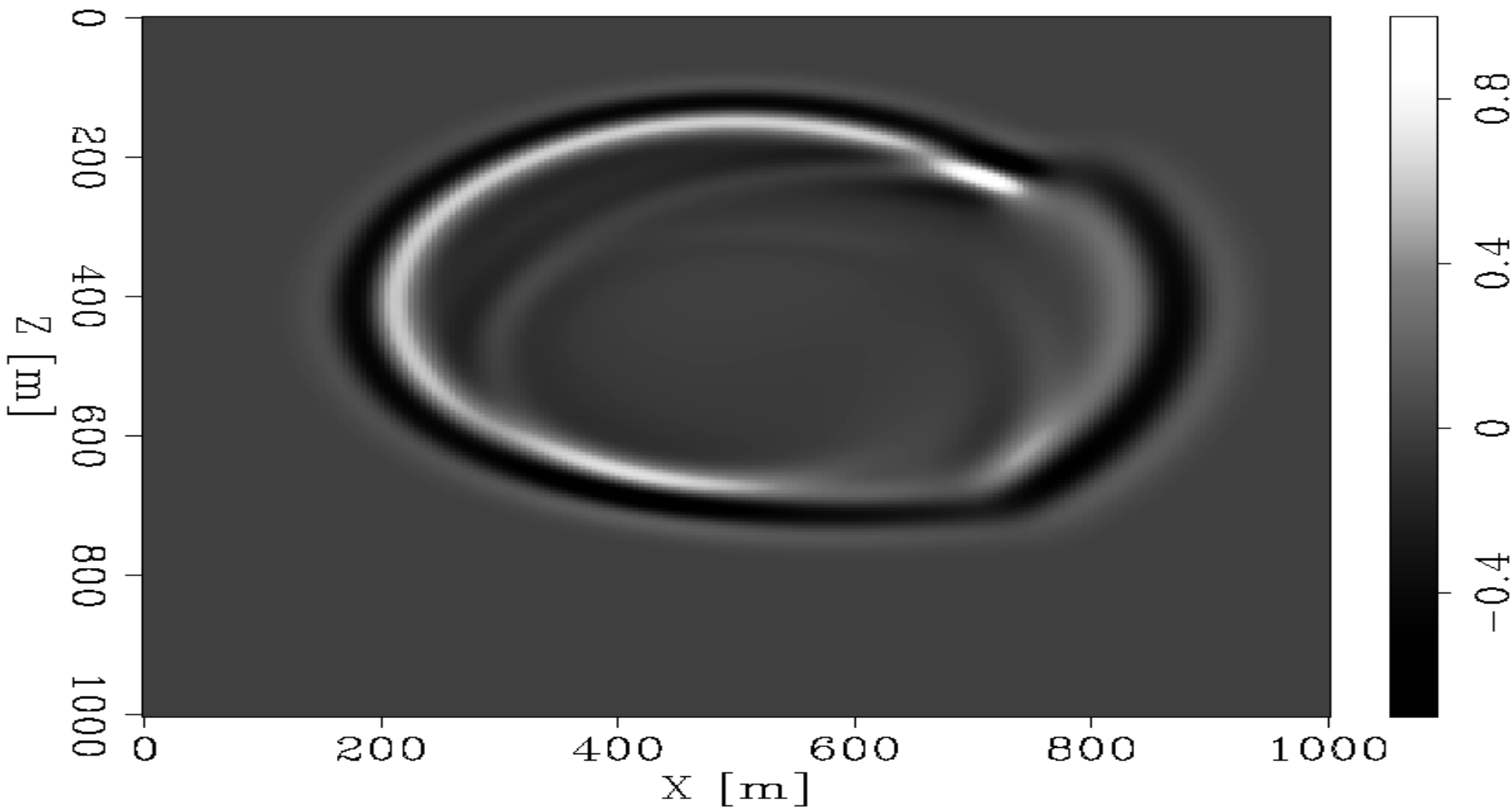
Error of 51 refs (100x)



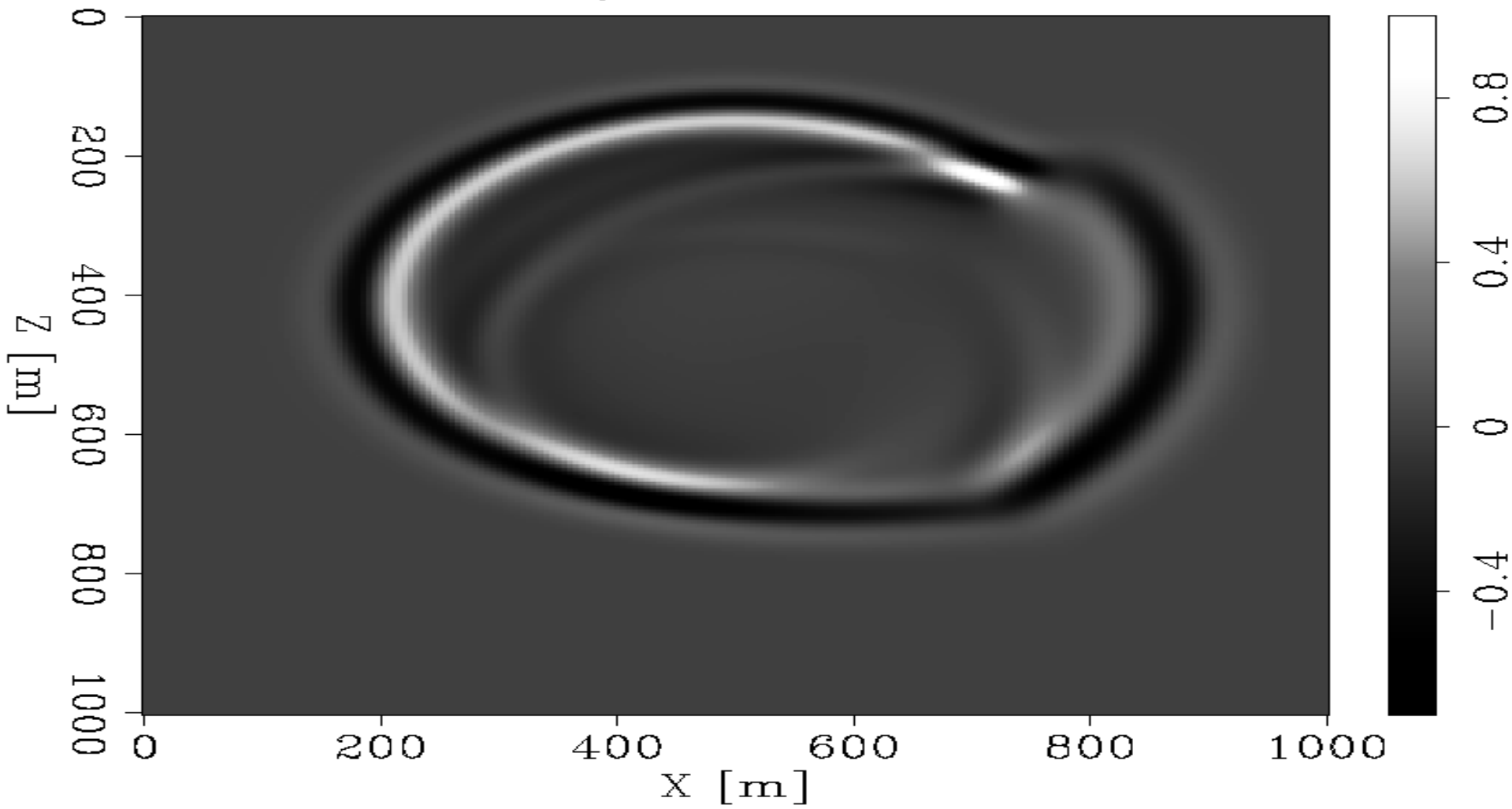
Operator interpolation: linear



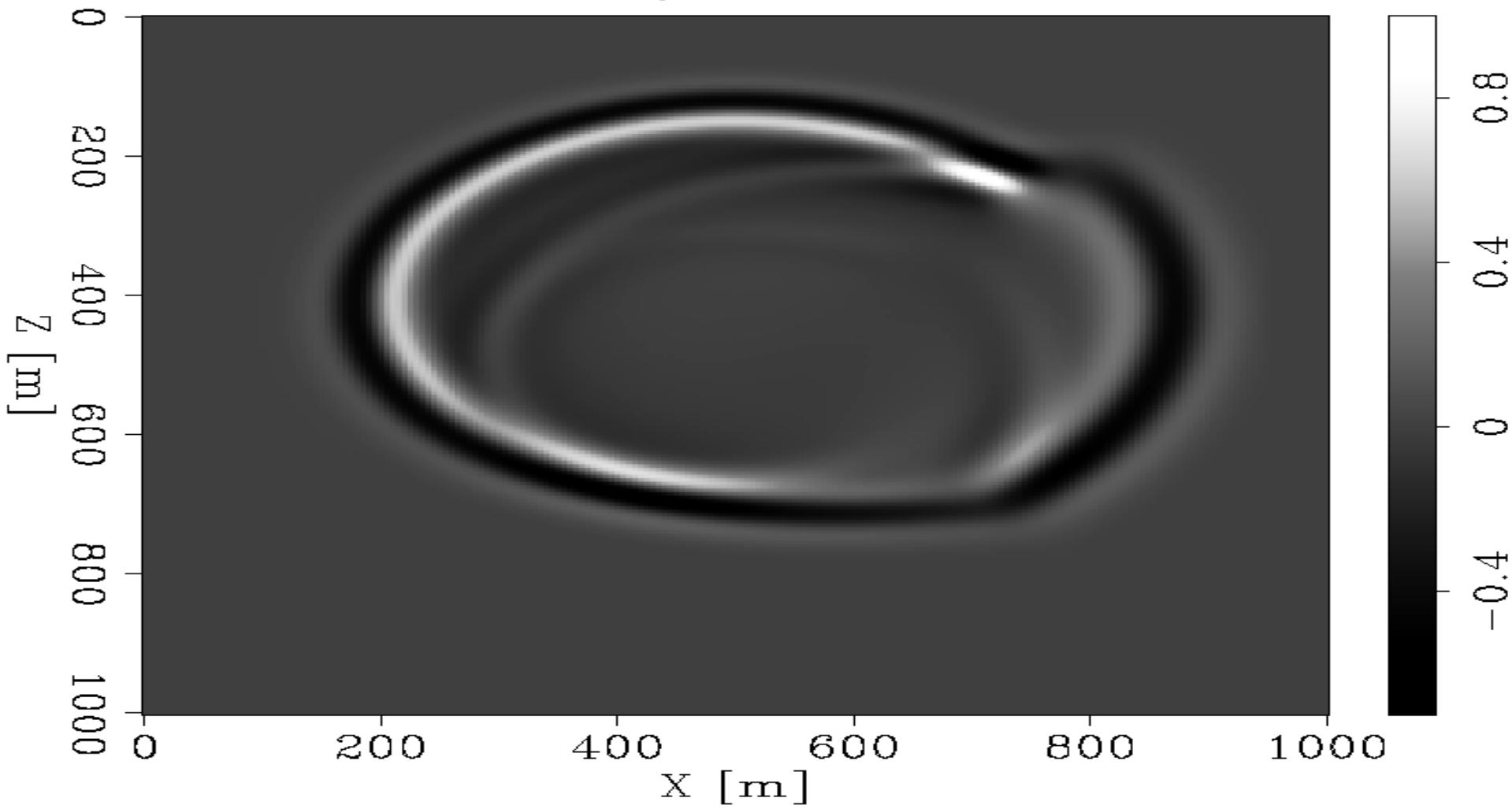
True wavefield



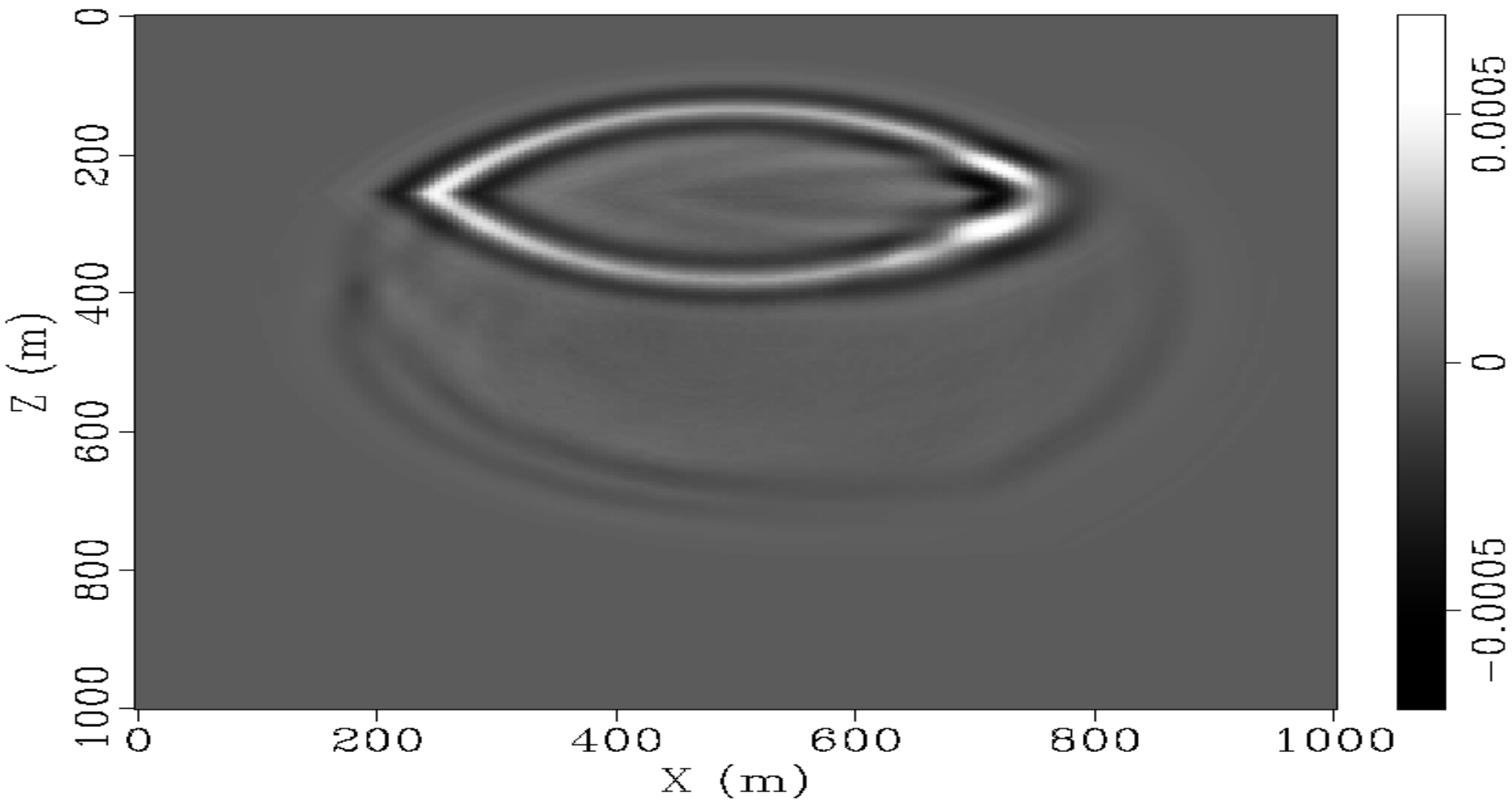
29 Lloyd refs, linear



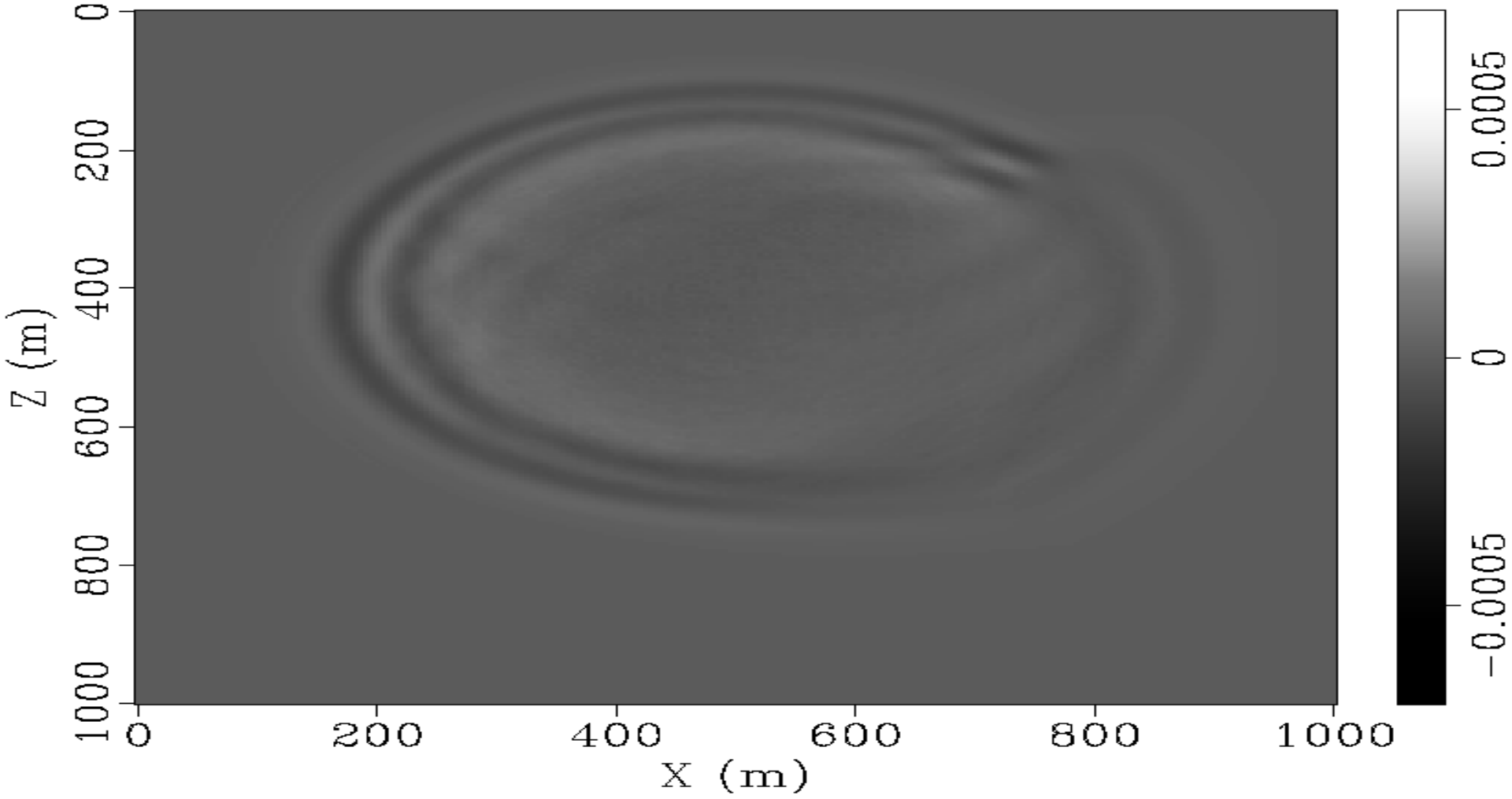
27 evenly refs, linear



Error of 29 Lloyd refs, linear (1000x)



Error of 27 evenly refs, linear (1000x)



Conclusions

- Shear artifacts can be removed completely by wavenumber-domain eigenvalue decomposition.
- Proposed method can be applied in wavenumber domain or space-time domain.
- Lower degree of anisotropy, smaller operator => possible hybrid scheme
- Dense operator is computationally expensive, but highly parallelizable and has better cache reuse.
- Application in inhomogeneous media by operator interpolation is stable and acceptably accurate.

Acknowledgements

- SEP sponsors
- Stanford School of Earth Sciences
- SEP colleagues, especially Ohad Barak and Elita (Yunyue) Li